

Electromagnetic (EM) Waves - ch. 34

Solution to Maxwell's equations in free space

$\Rightarrow$ Changing $\underline{E}$ field will induce a $\underline{B}$ field

We find solutions obey : $\frac{\partial^2 E}{\partial t^2} = c^2 \frac{\partial^2 E}{\partial x^2}$ (1)

with corresponding

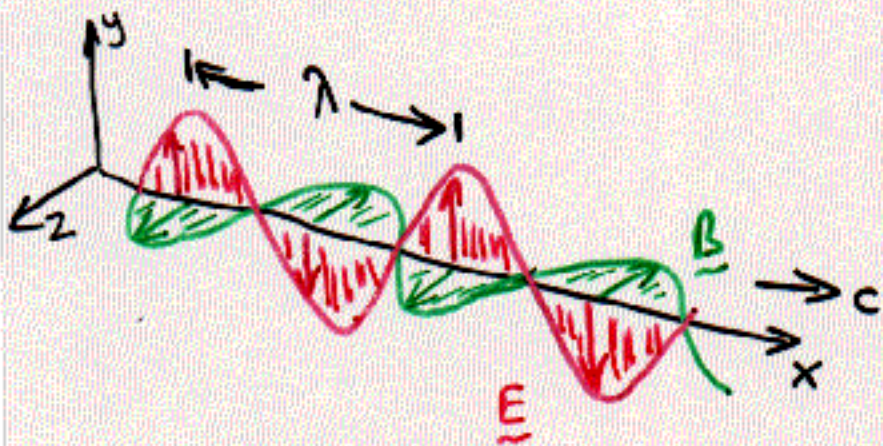
$$B = E/c \quad (2)$$

$$\text{Where } c = \frac{1}{\sqrt{\mu_0 \epsilon_0}}$$

Can verify that if $\underline{E}$ field is a wave : $\underline{E} = \underline{E}_0 \sin(kx - \omega t)$

(1) satisfied when speed of wave $\frac{\omega}{k} = c$: speed of light!

- $\underline{E}$, $\underline{B}$ are $\perp$ to each other, and direction of wave travel $= \underline{E} \times \underline{B}$ (transverse wave)
- $\underline{E}$, $\underline{B}$ are in phase with $B = E/c$



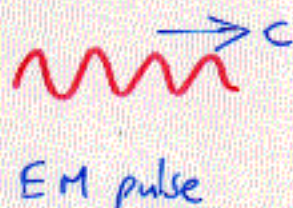
Linked $\underline{E}$, $\underline{B}$ fields
propagate together

EM Waves Transport energy and momentum in fields

- true even in vacuum



mid-flight:
energy, momentum in particle



mid-flight:
energy, momentum in $\underline{E}$, $\underline{B}$

$\underline{E}$, $\underline{B}$ fields exert a force on charged particles

$$\underline{F} = q (\underline{E} + \underline{v} \times \underline{B})$$

$\Rightarrow$ transfer of energy and momentum

(easier to measure than $\underline{E}$, $\underline{B}$ directly for light)

As for sound, Intensity $[W/m^2] \propto (\text{amplitude})^2$

With $\underline{E} = \underline{E}_0 \sin(kx - \omega t)$ and $\underline{B} = \underline{E}/c$ we find:

$$\underline{I} = \left(\frac{1}{\mu_0} \underline{E} \underline{B} \right)_{av} = \frac{1}{2} \cdot \frac{1}{c \mu_0} E_0^2 \quad [W/m^2]$$

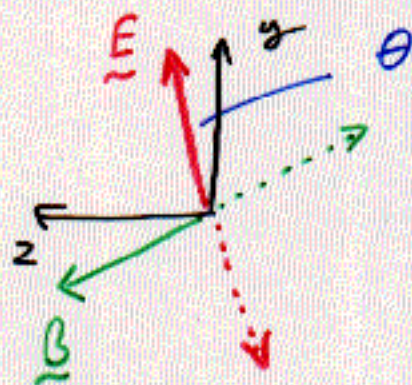
with electric field $\underline{E}$ in Volt/meter: $\underline{E} = \frac{\partial V}{\partial x}$

Polarization of EM Waves

$\underline{E}$, $\underline{B}$ fields are $\perp$ to direction of travel

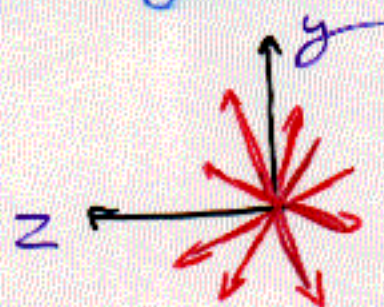
Plane containing $\underline{E}$ vector is plane of polarization

Viewed end-on :



Some sources produce $\sim 100\%$ polarized EM radiation
(e.g. TV and radio transmitters)

Most sources produce a jumble of EM waves, random phases



un-polarized light

Can represent $\underline{E}$ as two $\perp$ polarized waves with equal intensity, random phase:

$$\underline{E} = \underline{E}_y + \underline{E}_z$$



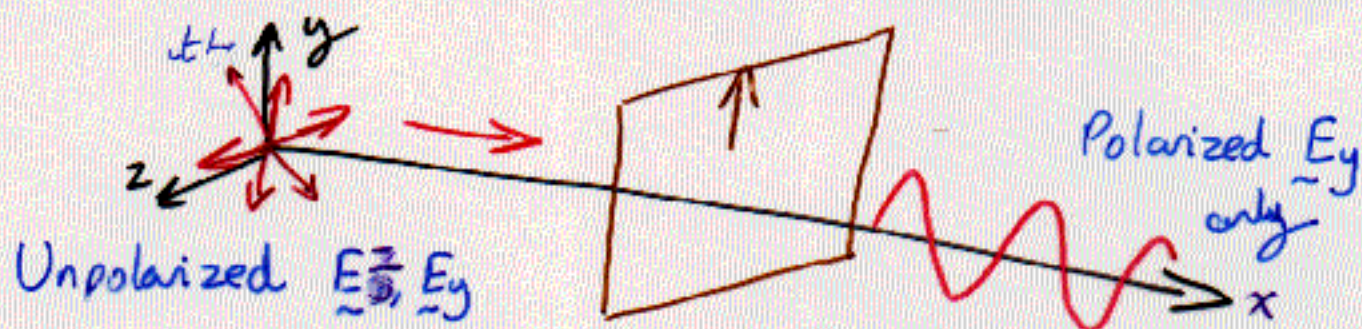
$$I_y = I_z = \frac{1}{2} I_0$$

↑
total intensity

Polarizing Filters

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Some materials (e.g. polaroid) only allow one plane of $\underline{E}$, $\underline{B}$ to be transmitted (+ absorb the rest)



$$I_0 = \frac{1}{2} \frac{1}{\mu_0 c} E_0^2$$

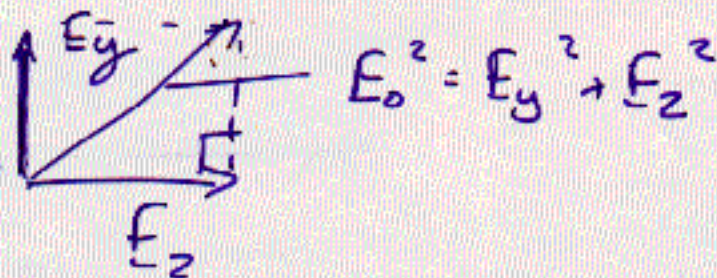
$$I_1 = \frac{1}{2} \frac{1}{\mu_0 c} E_y^2$$

$$I_0 = \frac{1}{2} \frac{1}{\mu_0 c} (E_x^2 + E_y^2) \quad \text{with, on average } E_x^2 = E_y^2$$

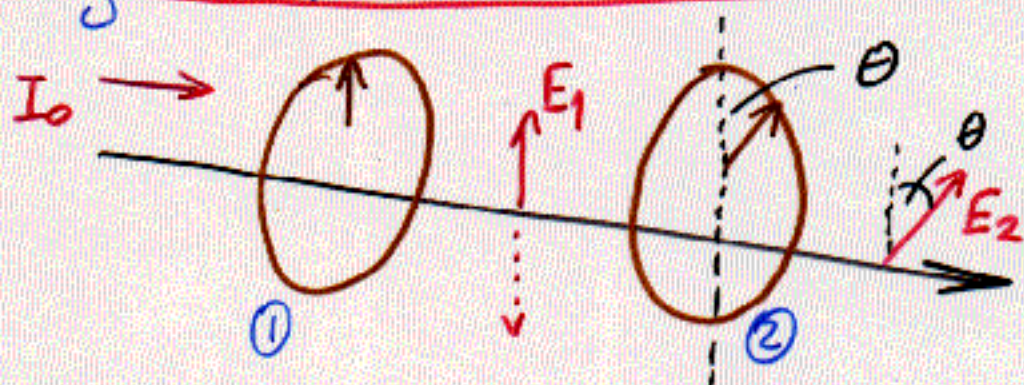
$$\Rightarrow \text{transmitted intensity } I_1 = \frac{1}{2} I_0$$

i.e. single polarizer cuts unpolarized intensity by a factor of 2

(Adding a 2nd polarizer with $\parallel$ transmission axis should make no difference: E_y transmitted 100%.)

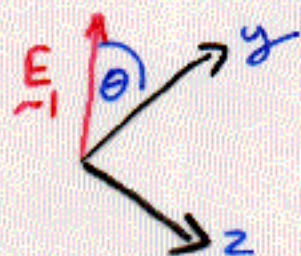


Combining two polarizers: Malus' Law



Take 2 polarizers, axes mis-aligned by θ

Unpolarized light incident on ① $\Rightarrow I_1 = \frac{1}{2} I_0$, with plane polarized beam. Take transmission axis of ② as the y-axis:



$$\underline{E}_1 = \underline{E}_y + \underline{E}_z$$

y-component of $\underline{E}_1$ is $E_1 \cos \theta$, with z-component absorbed
 $\Rightarrow$ transmitted $|\underline{E}_2| = |\underline{E}_y| = E_1 \cos \theta$

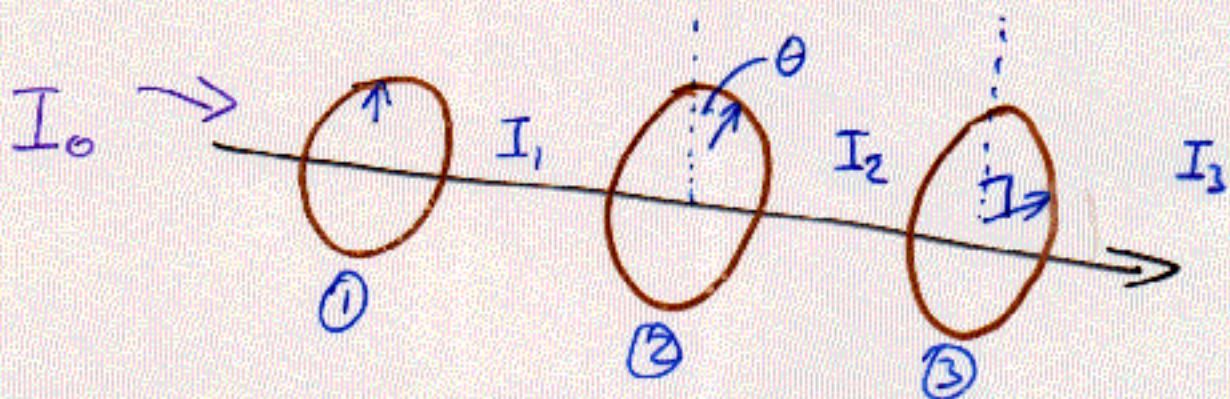
So since intensity $\propto E^2$

$$I_2 = I_1 \cos^2 \theta$$

— convenient means to vary intensity of a light source
 (from 0 to $\frac{1}{2} I_0$).

Application: Liquid Crystal Displays (LCDs)

Example: 3 Polarizers : first + last at 90°



We can rotate polarized light by 90° using 3 polaroid sheets

$I_2 = I_1 \cos^2 \theta$ as before, transmitted at θ to the vertical.

After polarizer ③: $I_3 = I_2 \cos^2 (90^\circ - \theta) \neq 0$ if $\theta \neq 0$

e.g. for unpolarized incident $I_0 = 800 \text{ W/m}^2$ and $\theta = 30^\circ$

After ①: $I_1 = \frac{1}{2} I_0 = \underline{400 \text{ W/m}^2}$

②: $I_2 = I_1 \cos^2 \theta = I_1 \left(\frac{\sqrt{3}}{2} \right)^2 = \frac{3}{4} I_1 = \frac{3}{8} I_0 = \underline{300 \text{ W/m}^2}$

③ $I_3 = I_2 \cos^2 60^\circ = I_2 \left(\frac{1}{2} \right)^2 = \frac{1}{4} I_2 = \frac{3}{16} I_0 = \underline{75 \text{ W/m}^2}$

Can show that I_3 is maximum when $\theta = 45^\circ$

(Then $I_3 = \frac{1}{4} I_0$)