

Rotational Equilibrium + Torque

21/1

e.g. Exert equal but opposite forces each side of steering wheel:



$$\vec{F}_L = -\vec{F}_R$$

No net force? $\checkmark$ So no net motion? $\times$

$$\vec{F}_L + \vec{F}_R = 0$$

Wheel starts to rotate, $\alpha = \frac{d\omega}{dt} \neq 0$

(though wheel does not translate)

Reason? $\vec{F}_L, \vec{F}_R$ do not act at the same point of application.

$\therefore$ Both $\vec{F}_L, \vec{F}_R$ provide a TORQUE ("twist") which does not cancel.

Define Torque (or moment) of force about
a rotation center as:
(pivot)

$$\tau = r F_{\perp} = r F \sin \theta$$

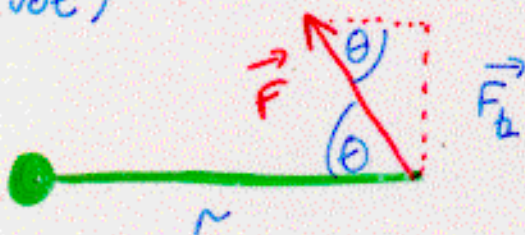
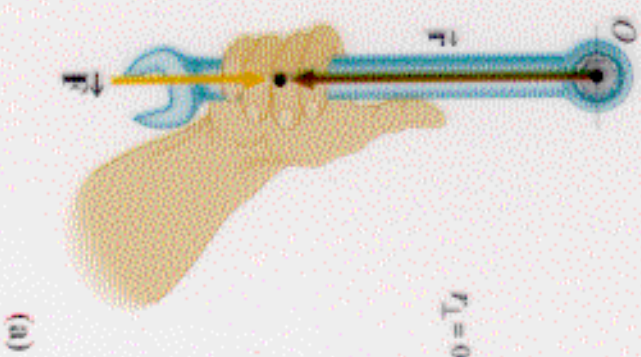


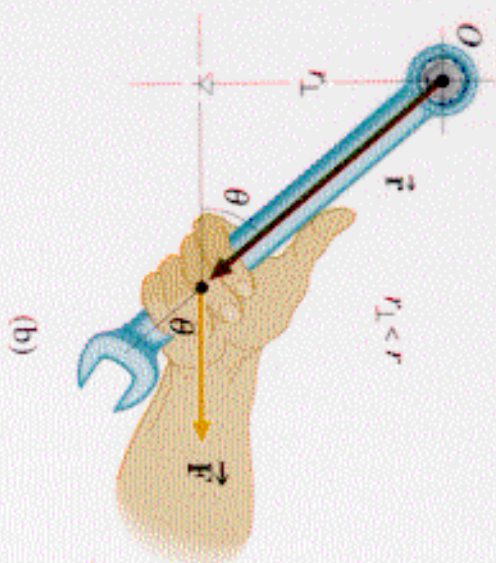
Figure 8.14

Torque =

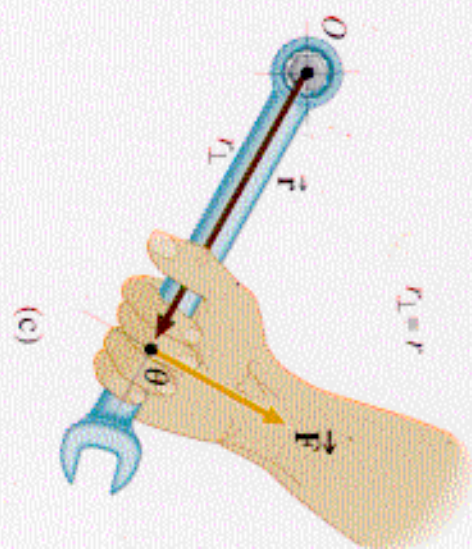
$$rF\sin\theta = r_{\perp}F\sin\theta = rF_{\perp}\sin\theta$$



$$\tau = 0$$



$$\tau = rF\sin\theta$$



$$\tau = rF$$

$$(\sin\theta = 1)$$

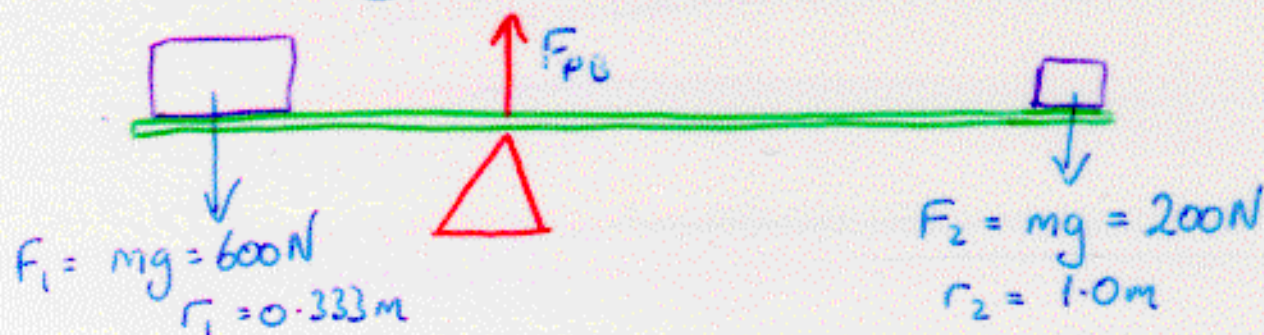
Rotational Eq. m. cont/d.

21/3

So if $\sum \vec{F} = \vec{F}_1 + \vec{F}_2 + \dots = 0$: no linear acceleration

But only if $\sum \tau = F_1 r_1 \sin \theta_1 + F_2 r_2 \sin \theta_2 + \dots = 0$ *
do we get no angular acceleration.

E.g. Balance weights, beam on pivot.



For "balance" ($\omega = 0$, $\alpha = \frac{d\omega}{dt} = 0$)

choose anti-clockwise as positive direction for ω

Then torques $F_1 r_1 = -F_2 r_2 = 600 \times 0.333\text{m} = 200\text{Nm}$
 $F_1 r_1 + F_2 r_2 = 0$

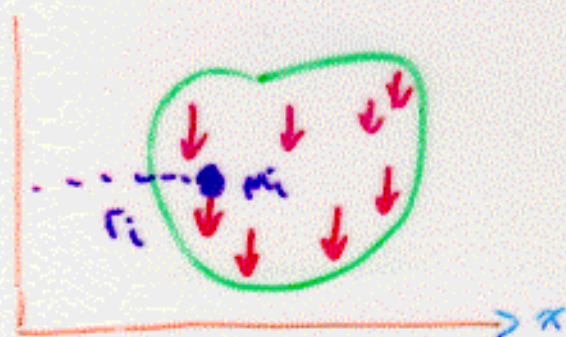
Note: $F_1 + F_2 = 600 + 200\text{N} \neq 0$, but additional

force of pivot on beam $F_{PB} = 800\text{N}$ in opposite direction

(and since $r = 0$, this force exerts no torque)

$\Rightarrow$ no net translation OR rotation.

Center of Gravity Extended Objects



Hold irregular object
in outstretched hand

To calculate torque, can add up $\sum m_i r_i$

Easier to define center of gravity : point
through which total weight can be thought to act.

i.e. Torque $\tau = \sum m_i g x_i \equiv M g x_c$ where $M = \sum m_i$

$$\text{i.e. } x_c = \frac{\sum m_i x_i}{\sum m_i}$$

Similarly for y_c, z_c in 3-dimensions



Center of Gravity and Stability

1. Hang object from one edge and let go:



$$\tau = r_b F \neq 0$$

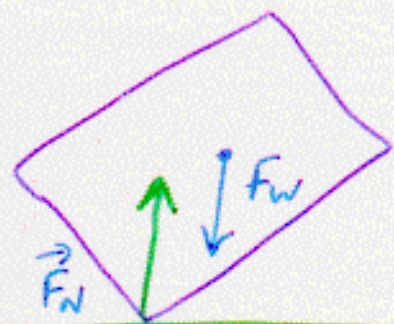
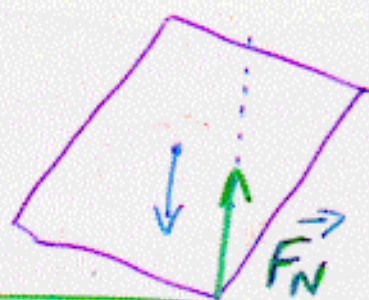


Net torque rotates object until $r \perp F$ and $\tau = 0$

$\Rightarrow$ easy method to locate c.g. of object

Note: C.g. not always located inside object,
e.g. Crescent or banana shape

2. Rest object on horizontal surface, then try to tip over (to right):



As long as c.g. stays to left of pivot, torque

$\tau = r_b F_w$ acts to restore equilibrium

($\vec{F}_N = -\vec{F}_w$ but exerts no torque about pivot)

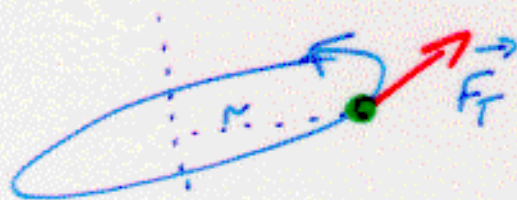
Rotational Dynamics

21/6

Already have 1st law of R.D:

- ① A body will remain at rest or rotate with $\omega = \text{constant}$ unless acted on by external torque

Now go back to small mass moving in circle



Tangential accel. $a_t = r \frac{d\omega}{dt} = r\alpha$ from before

Newton II : $F_t = ma_t = mr\alpha$

This force exerts a torque with F_t & r

$$\tau = F_t \cdot r = mr^2\alpha$$

If we define Moment of Inertia $I = mr^2$

② $\tau = I\alpha$ (c.f. $F = ma$)

For many point masses in a rigid body (same ω , same α)

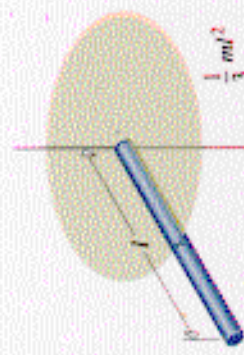
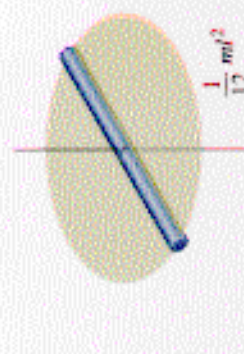
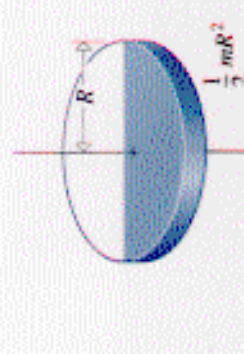
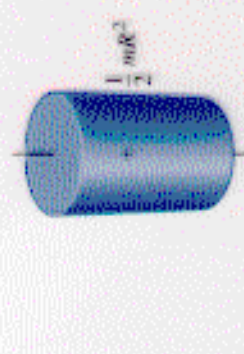
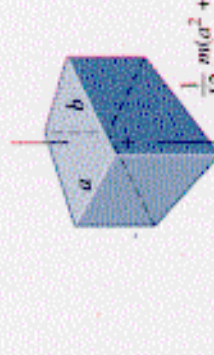
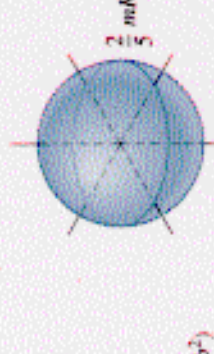
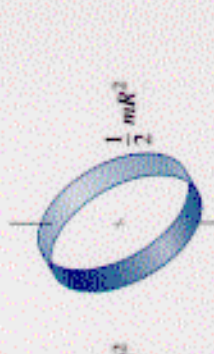
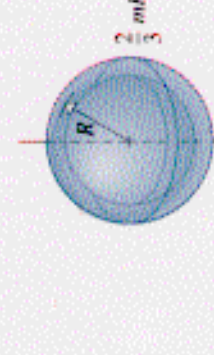
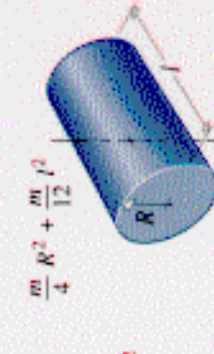
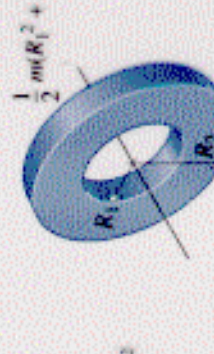


Total torque

$$\tau = \sum m_i r_i^2 \alpha = I\alpha$$

where $I = \sum m_i r_i^2 = \int r^2 \cdot dm$

Moments of inertia

 Slender rod (a) $\frac{1}{3} ml^2$	 Slender rod (b) $\frac{1}{12} ml^2$	 Disk (c) $\frac{1}{2} mR^2$	 Cylinder (d) $\frac{1}{2} mR^2$
 Rectangular parallelepiped (e) $\frac{1}{12} m(a^2 + b^2)$	 Solid sphere (f) $\frac{2}{5} mR^2$	 Hoop (g) mR^2	 Hoop (h) $\frac{1}{2} mR^2$
 Spherical shell (h) $\frac{2}{3} mR^2$	 Cylinder (i) $\frac{m}{4} R^2 + \frac{m}{12} l^2$	 Cone (j) $\frac{3}{10} mR^2$	 Hoop (k) $\frac{1}{2} m(R_1^2 + R_2^2)$

$$I = \int r^2 dm = \int r^2 \rho dV$$

depends on location and direction of rotation axis