

# Theory of Fine-Scale Zonal Flow Generation

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Lu Wang,\* T.S. Hahm

Princeton Plasma Physics Laboratory

\*Peking University, Beijing, China

Ackn: Y. Xiao, and Z. Lin

GPS-TTBP-Fest, APS-DPP, 2008

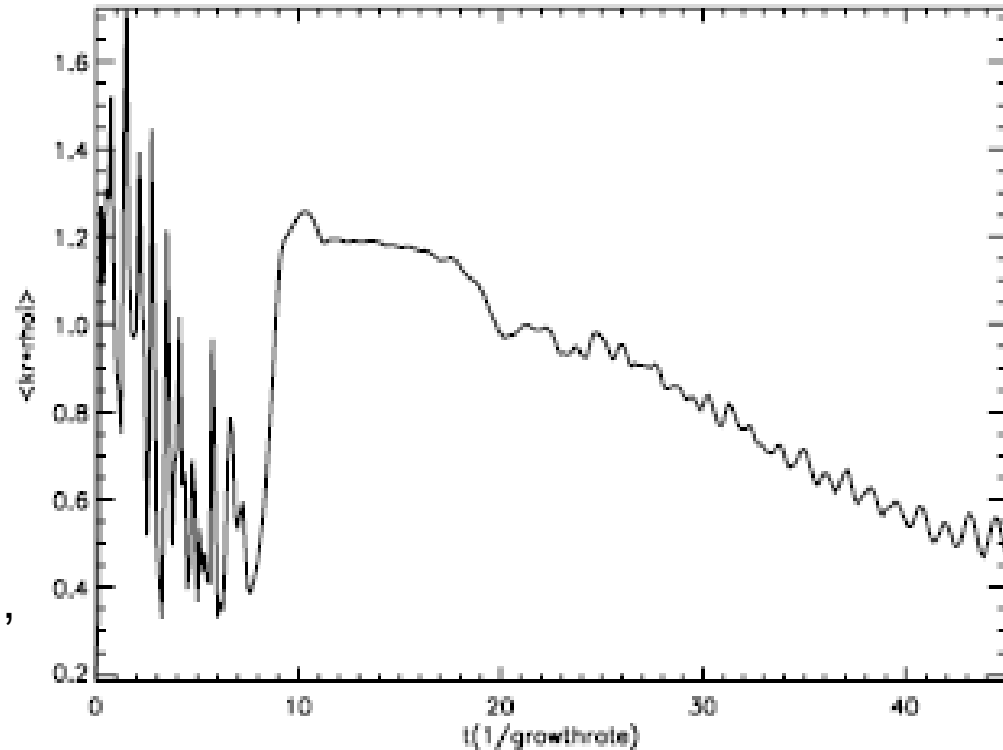
# Motivation

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- The RH neoclassical polarization  $1.6\varepsilon^{3/2}q_r^2\rho_{\theta i}^2$  [Rosenbluth and Hinton, PRL 80, 724 (1998)] was used in most ZF generation theories [P. H. Diamond et al., IAEA-CN-69/TH3/1 (1998) & L. Chen et al., PoP 7, 3129 (2000)].
- RH neoclassical polarization expression is valid for  $q_r\rho_{\theta i} \ll 1$  .

# Motivation

- A relatively **short scale** ZF ( $q_r \rho_{\theta i} \sim 1$ ) has been found in Nonlinear GTC simulation of Trapped Electron Mode (TEM) turbulence [Z. Lin et al., IAEA-CN/TH/P2-8 (2006) & Y. Xiao, et al., TTF Meeting (2008)].



This figure of Zonal flow length scale for CTEM is from Yong Xiao's work

# Turbulence driven radial current shielded by ZF

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$$en_0\chi_{total} \frac{\partial}{\partial t} \left( \frac{e\phi_z}{T_e} \right) + \frac{\partial}{\partial r} \left( \frac{\delta \langle J_r \rangle}{\delta \phi_z} \phi_z \right) = 0 \quad [\text{Diamond et al., IAEA (1998)}]$$

$\phi_z$  zonal flow potential,  $\langle J_r \rangle$  transport-induced radial current

$\chi_{total}$  polarization including neoclassical enhancement

- Poynting theorem for drift wave turbulence

$$J_r = -\frac{n_0 T_e}{2} \sum_k \frac{ck_\theta}{B} \frac{\partial}{\partial r} \left( \frac{\partial \chi}{\partial k_r} \Big|_{\omega_k} \left| \frac{e\phi_k}{T_e} \right|^2 \right), \quad \chi: \text{susceptibility}$$

[Diamond and Kim, PFB 3, 1626 (1991)]

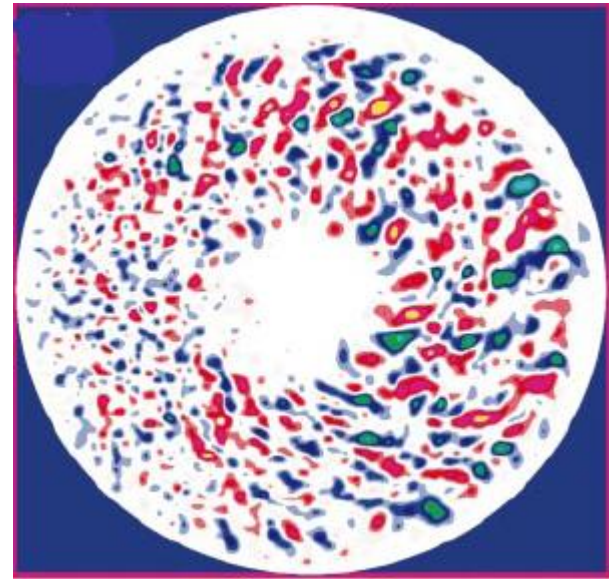
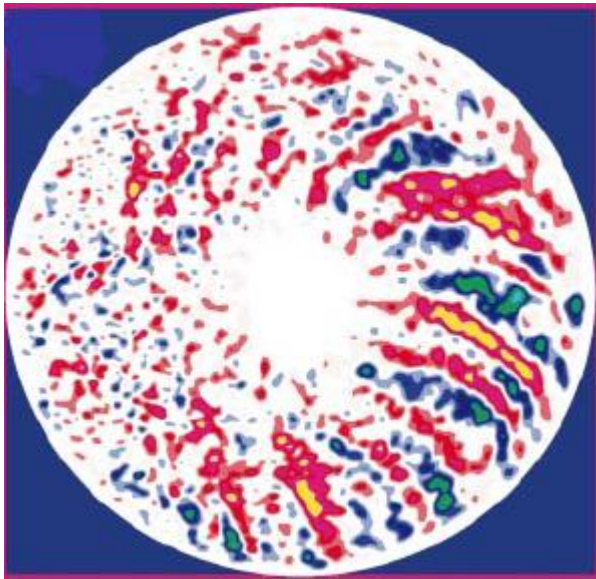
- Use the definition of wave action density N

[Lebedev et al., PoP 2, 4420 (1995)]

$$N = \frac{E_k}{\omega_k} = -\frac{n_0 T_e}{2} \frac{\partial \chi}{\partial \omega} \left| \frac{e\phi_k}{T_e} \right|^2 \quad \text{in wave kinetic approach}$$

# ZF Generation and Random Shearing Occur Simultaneously

- $\omega_k \gg \Omega_z \rightarrow$  Drift Wave action density,  $N$ , is conserved (adiabatic invariant)
- Random Shearing by Zonal Flows results in increase of  $k_r$  of DWT. [Diamond et al., IAEA '98 & Hahm et al., PoP '99]
- With  $E_k = N\omega_k$ , DW Energy decreases as  $\omega_k$  is down-shifted.
- $\frac{d}{dt}(E_{DW} + E_{ZF}) = 0, \rightarrow$  ZF growth [Diamond et al., IAEA '98]



# ZF Growth Rate from Wave Kinetic Approach

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- Wave kinetic equation [Diamond et al., IAEA (1998)]

$$\frac{\partial}{\partial t} N + (\vec{v}_g + \vec{v}_{ZF}) \cdot \frac{\partial}{\partial \vec{x}} N - \frac{\partial}{\partial \vec{x}} (\omega + \vec{k} \cdot \vec{v}_{ZF}) \cdot \frac{\partial}{\partial \vec{k}} N = \gamma_k N - \frac{\omega_k}{N_0} N^2$$

- Using scale separation and action response to ZF variation,  $\tilde{N}$ , we obtain the ZF growth rate

$$\gamma_z = -\frac{1}{\chi_{total}} q_r^4 c_s^2 \rho_s^2 \sum_k k_\theta^2 \frac{\partial \chi}{\partial k_r} \Big|_{\omega_k} \frac{2 \partial \langle N \rangle / \partial k_r}{n_0 T_e \partial \chi / \partial \omega} R(k, q_r) \left( 1 - \frac{\chi_z}{\omega_{*e} / \omega_k} \right)$$

Expect growth  $\gamma_z > 0$  for  $\partial \langle N \rangle / \partial k_r < 0$

# Neo-polarization for $q_r \rho_{\theta i} \sim 1$

- Neo-Polarization from bounce-kinetic equation (via action-angle variables  $\rightarrow$  conventional guiding center variables) [Fong and Hahm, PoP 6, 188 (1996)]

- For  $q_r \rho_{\theta i} \ll 1$  :  $\chi_{nc}^{tr} = 1.35 \tau \varepsilon^{3/2} q_r^2 \rho_{\theta i}^2$

The same scaling with RH expression

- For  $q_r \rho_{\theta i} \gg 1$  :  $\chi_{nc}^{tr} = 0.9 \tau \sqrt{\varepsilon} \left(1 - \frac{0.2}{\sqrt{\varepsilon} q_r \rho_{\theta i}}\right)$

- The generalized analytic expression

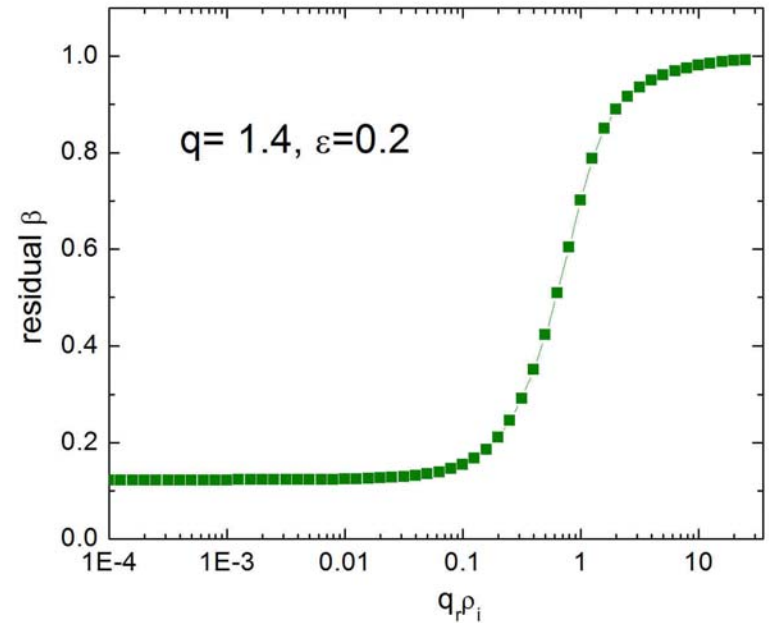
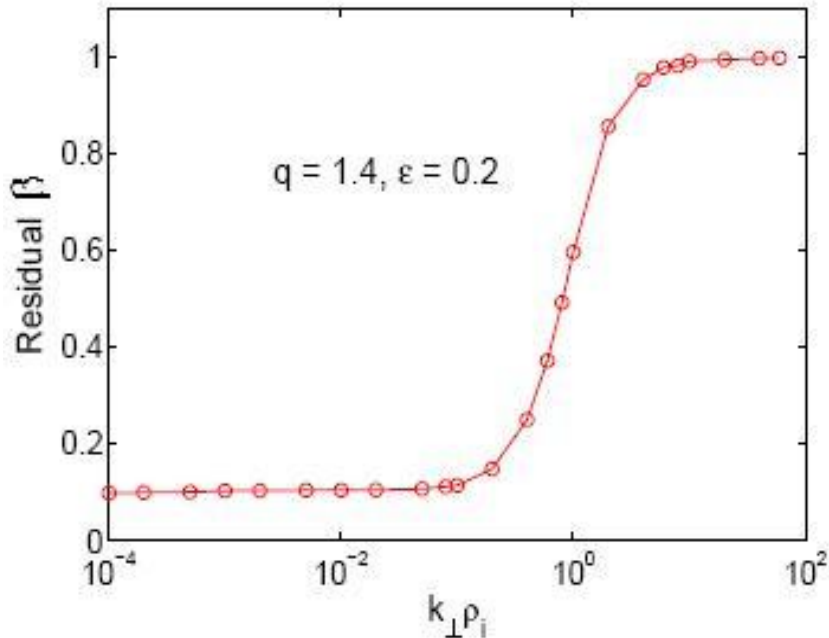
$$\chi_{total} = \frac{\tau}{\frac{1}{1.6 \varepsilon^{3/2} q_r^2 \rho_{\theta i}^2} + \frac{1}{0.9 \sqrt{\varepsilon} \left(1 + \frac{0.2}{\sqrt{\varepsilon} q_r \rho_{\theta i}}\right)}} + \tau (1 - 0.9 \sqrt{\varepsilon}) \left(1 - \Gamma_0(\varepsilon q_r \rho_{\theta i})\right)$$

 barely circulating particles

 strongly circulating particles

# Results (Neo-polarization)

For arbitrary  $q_r \rho_{\theta i}$  Residual zonal flow  $\beta = \chi_{cl} / \chi_{total}$

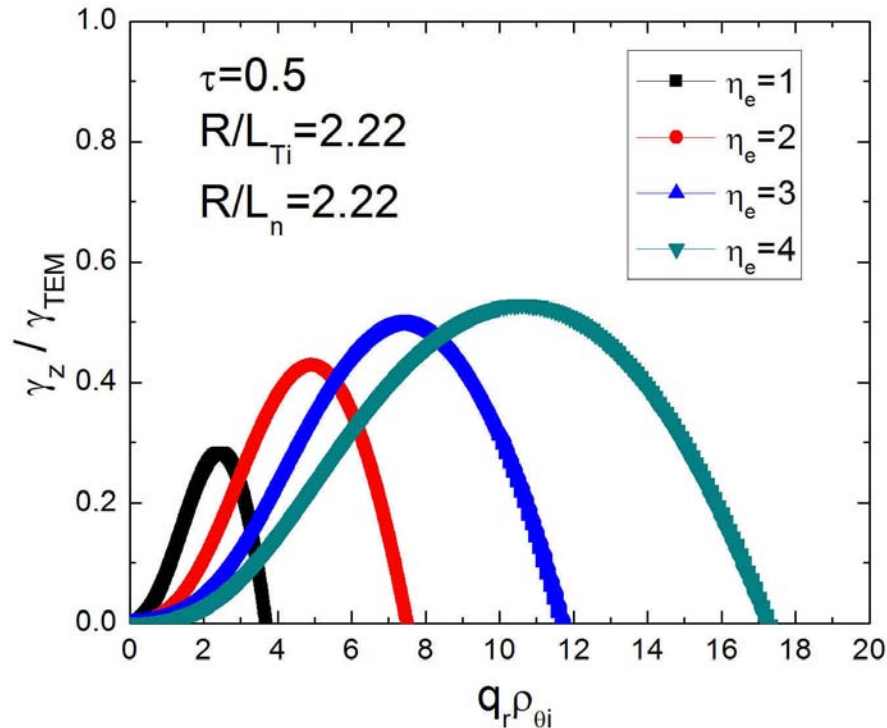


The left figure is taken from Y. Xiao's PhD thesis(2006): Fig 4-5 (numerical)  
The right one is our generalized analytic result.

Our analytic expression agrees with  
Y. Xiao's numerical result very well!



# ZF Growth Rate for TEM



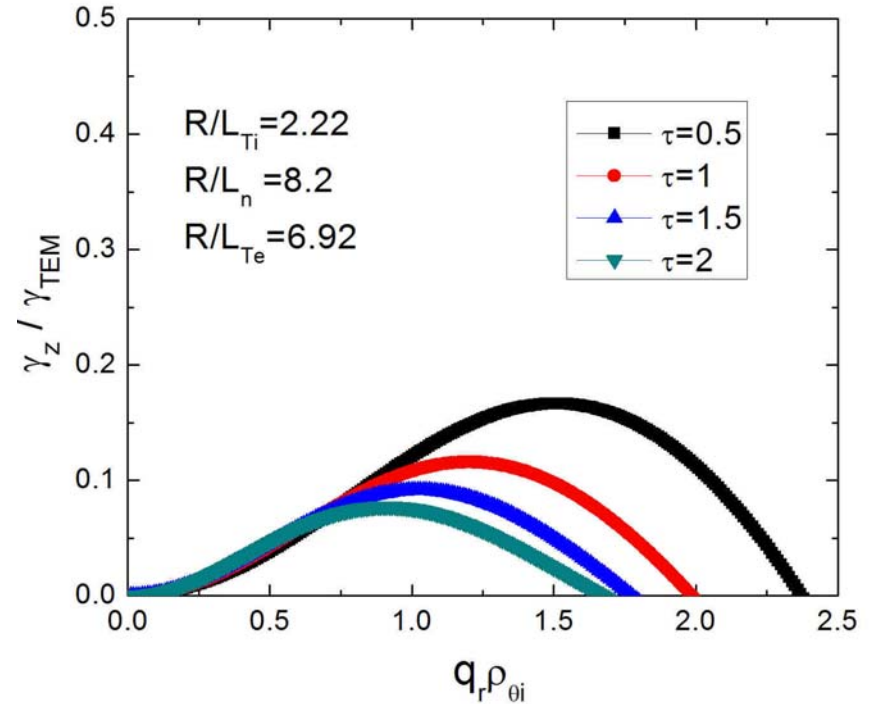
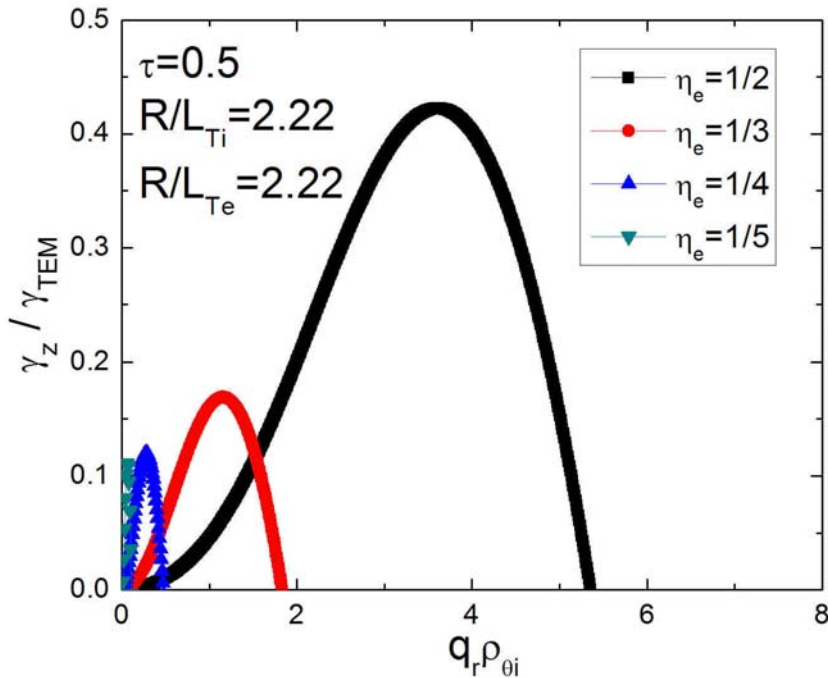
Using the susceptibility & linear growth rate for CTEM and the generalized polarization

This figure is for  $k_{\perp} \rho_i \sim 0.2$

and  $\varepsilon = 0.18$

The spatial scale of ZF at which  $\gamma_z / \gamma_{TEM}$  maximizes  $\sim$  ZF scale from Z. Lin (2006) and Y. Xiao (2008)!  
Normalized ZF growth rate and characteristic  $q_r$  increase with an increase in electron temperature gradient.

# ZF Growth Rate for TEM



These figures are for  $k_{\perp} \rho_i \sim 0.2$  and  $\varepsilon = 0.18$

Normalized ZF growth rate and characteristic  $q_r$  increase with a decrease in density gradient and temperature ratio  $T_e/T_i$ .

# Conclusion

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- Generalized neoclassical polarization expression for full banana width  $q_r \rho_{\theta i} \sim 1$  is obtained from bounce-kinetic equation.
- ZF Growth for short spatial scale is obtained.
- Qualitatively consistent with GTC simulation results (Yong's Job)
- Open questions:
  - ZF saturation mechanism still needs to be addressed.
  - Role of trapped electron nonlinearity [previously concluded to be small in Lin et al., IAEA (2004)]
  - Robustness of WKE approach for short scale ZF

**Thank you!**