

Properties of dark matter haloes and their correlations: the lesson from principal component analysis

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ABSTRACT

We study the correlations between the structural parameters of dark matter haloes using principal component analysis. We consider a set of eight parameters, six of which are commonly used to characterize dark matter halo properties: mass, concentration, spin, shape, overdensity and the angle (Φ_L) between the major axis and the angular momentum vector. Two additional parameters (x_{off} and ρ_{rms}) are used to describe the degree of ‘relaxedness’ of the halo. We find that we can account for much of the variance of these properties with halo mass and concentration, on the one hand, and halo relaxedness on the other. None the less, three principal components are usually required to account for most of the variance. We argue that halo mass is not as dominant as expected, which is a challenge for halo occupation models and semi-analytic models that assume that mass determines other halo (and galaxy) properties. In addition, we find that the angle Φ_L is not significantly correlated with other halo parameters, which may present a difficulty for models in which galaxy discs are oriented in haloes in a particular way. Finally, at fixed mass, we find that a halo’s environment (quantified by the large-scale overdensity) is relatively unimportant.

Key words: gravitation – methods: numerical – methods: statistical – galaxies: haloes – cosmology: theory – dark matter.

1 INTRODUCTION

In the paradigm of hierarchical structure formation, gravitational evolution causes dark matter (DM) particles to cluster around peaks of the initial density field and to collapse into virialized objects (haloes), which provide the potential wells in which galaxies subsequently form (White & Rees 1978). It is therefore expected that the properties of a galaxy are correlated with the properties of its host halo. Small haloes merge to form larger and more massive haloes, which tend to be located in dense environments and are expected to host groups of galaxies.

It is commonly assumed that the mass of a DM halo determines how the galaxies it hosts form and evolve; other halo properties are assumed to be less important at fixed mass. This assumption is usually made both by those who use halo models of galaxy clustering (e.g. Zehavi et al. 2005; Skibba et al. 2006; van den Bosch et al. 2007; Moster et al. 2010) and by those who use semi-analytic models of galaxy formation (e.g. Bower et al. 2006; Cattaneo et al. 2007; De Lucia & Blaizot 2007; Somerville et al. 2008). Many of such models have recently been able to provide a description, and perhaps a plausible explanation, of various galaxy statistics, such

as the luminosity function, stellar mass function, colour–magnitude distribution, correlation function and mark correlation functions, lending credence to claims about the importance of halo mass.

None the less, although halo properties, such as their mass, spin, concentration and shape, are correlated with each other, there is considerable scatter in these correlations (e.g. Avila-Reese et al. 2005; Macciò et al. 2007; Macciò, Dutton & van den Bosch 2008; Ragone-Figueroa et al. 2010). In addition, recent studies of numerical simulations have shown that halo formation time is correlated with the environment at fixed mass (e.g. Sheth & Tormen 2004; Gao, Springel & White 2005; Dalal et al. 2008). Some have argued that the halo occupation distribution (used in models of galaxy clustering) is correlated with formation time at fixed mass (Zentner et al. 2005a; Wechsler et al. 2006; Giocoli et al. 2010). It has also been argued that such halo ‘assembly bias’ affects the clustering of galaxies and that the effect varies with luminosity and colour (Croton, Gao & White 2007). None the less, careful studies of the environmental dependence of galaxy luminosity and colour in the real Universe have explained this dependence in terms of the environmental dependence of halo mass alone (e.g. Tinker et al. 2008a; Skibba & Sheth 2009). Such models are also encouraged by an analysis of a halo-based galaxy group catalogue, in which, given a halo’s mass, the colours of the galaxies it hosts are not significantly correlated with the large-scale density field (Blanton & Berlind 2007).

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In addition, the halo model can reproduce the observed non-monotonic relation between clustering strength and density (Abbas & Sheth 2007); note, however, that assembly bias effects are thought to be stronger for lower mass haloes (e.g. Croton et al. 2007; Wang, Mo & Jing 2007), which host fainter galaxies than the Sloan Digital Sky Survey can reliably probe. The results of these studies suggest that assembly bias affects the formation of observed galaxies only weakly, if at all, although the debate is not yet resolved.

It is also commonly assumed that typical DM haloes are relaxed and in virial equilibrium. Nevertheless, Macciò et al. (2007) have shown that many haloes are significantly unrelaxed, sometimes because they are experiencing or have recently experienced a merger or have not yet virialized. In addition, statistical studies of galaxy groups and clusters have shown that the most massive galaxy in these systems is often significantly offset from the centre of the potential well, which may be an indication that some of the systems are unrelaxed (van den Bosch et al. 2005; Skibba et al. 2011).

The relation between halo formation and galaxy formation is far from clear, and these issues are in need of further investigation. For example, to what extent halo mass determines other halo properties, and to what extent relaxed and unrelaxed haloes differ, remains to be seen. The purpose of this paper is to examine the properties of DM haloes in numerical simulations, with principal component analysis (PCA). This type of analysis allows us to quantify the correlations between the halo properties, determine which properties are the most important, and to what extent halo mass is the dominant property. It also allows us to characterize the distinctive properties of relaxed and unrelaxed haloes.

This paper is organized as follows. We describe the DM halo simulations in Section 2. In Section 3, we describe our PCAs and our choice of halo properties to analyse. We then present the PCA results in Section 4. We further examine correlations with halo mass and relaxedness in Section 5. Finally, we end with a brief discussion of our results and their implications in Section 6.

2 NUMERICAL SIMULATIONS

Table 1 lists all of the simulations used in this work. Most of them have been already presented in Macciò et al. (2008) and Muñoz-Cuartas et al. (2011). We have run simulations for several different box sizes, which allows us to probe halo masses covering the entire range $10^{10} \lesssim M \lesssim 10^{15} h^{-1} M_{\odot}$. In addition, in some cases we have run multiple simulations for the same cosmology and box size, in order to test for the impact of cosmic variance (and to increase the final number of DM haloes).

All simulations have been performed with PKDGRAV, a tree code written by Joachim Stadel and Thomas Quinn (Stadel 2001). The code uses spline kernel softening, for which the forces become

completely Newtonian at twice the softening length, ϵ . The physical values of ϵ at $z = 0$ are listed in Table 1. The initial conditions are generated with the GRAFIC2 package (Bertschinger 2001).

We have set the cosmological parameters according to the fifth-year results of the *Wilkinson Microwave Anisotropy Probe* mission (Kogut et al. 2009), namely $\Omega_m = 0.258$, $\Omega_L = 0.742$, $n = 0.963$, $h = 0.72$ and $\sigma_8 = 0.796$.

In all of the simulations, DM haloes are identified using a spherical overdensity (SO) algorithm. We use a time-varying virial density contrast determined using the fitting formula presented in Mainini et al. (2003). We include in the halo catalogue all the haloes with more than 1000 particles (see Macciò et al. 2008, for further details on our halo-finding algorithm). Note that we include only host DM haloes in our analysis; halo substructures are not included.

2.1 Halo parameters

For each SO halo in our sample we determine a set of parameters, including the virial mass and radius, the concentration parameter, the angular momentum, the spin parameter and axial ratios (shape). Below, we briefly describe how these parameters are defined and determined. A more detailed discussion can be found in Macciò et al. (2007, 2008).

2.1.1 Concentration parameter

To compute the concentration of a halo, we first determine its density profile. The halo centre is defined as the location of the most bound halo particle, and we compute the density (ρ_i) in 50 spherical shells, spaced equally in logarithmic radius. Errors on the density are computed from the Poisson noise due to the finite number of particles in each mass shell. The resulting density profile is fitted with a Navarro–Frenk–White (NFW) profile (Navarro, Frenk & White 1997):

$$\frac{\rho(r)}{\rho_{\text{crit}}} = \frac{\delta_c}{(r/r_s)(1+r/r_s)^2}. \quad (1)$$

During the fitting procedure we treat both r_s and δ_c as free parameters. Their values, and associated uncertainties, are obtained via a χ^2 minimization procedure using the Levenberg–Marquardt method. We define the rms of the fit as

$$\rho_{\text{rms}} = \frac{1}{N} \sum_i^N (\ln \rho_i - \ln \rho_m)^2, \quad (2)$$

where ρ_m is the fitted NFW density distribution. Finally, we define the concentration of the halo, $c_{\text{vir}} \equiv R_{\text{vir}}/r_s$, using the virial radius obtained from the SO algorithm, and we define the error on $\log c$ as $(\sigma_{r_s}/r_s)/\ln(10)$, where σ_{r_s} is the fitting uncertainty on r_s .

2.1.2 Spin parameter

The spin parameter is a dimensionless measure of the amount of rotation of a DM halo. We use the definition introduced by Bullock et al. (2001):

$$\lambda = \frac{J_{\text{vir}}}{\sqrt{2} M_{\text{vir}} V_{\text{vir}} R_{\text{vir}}}, \quad (3)$$

where M_{vir} , J_{vir} and V_{vir} are the mass, total angular momentum and circular velocity at the virial radius, respectively. See Macciò et al. (2007) for a detailed discussion and for a comparison of the different definitions of the spin parameter. Note that the concentration and spin parameters are not well defined in unrelaxed haloes; for this

Table 1. *N*-body simulation parameters.

Name	Box size (Mpc)	<i>N</i>	Part. mass ($h^{-1} M_{\odot}$)	Force soft. (h^{-1} kpc)	<i>N</i> _{halo} >500
W5-20.1	20	250 ³	1.37e7	0.43	974
W5-30.1	30	300 ³	2.67e7	0.64	984
W5-40.1	40	250 ³	1.09e8	0.85	1119
W5-90.1	90	300 ³	7.21e8	1.92	1998
W5-90.2	90	600 ³	9.02e7	0.85	16161
W5-180	180	300 ³	5.77e9	3.83	2302
W5-300.1	300	400 ³	1.13e10	4.74	5845
W5-300.2	300	400 ³	1.13e10	4.74	5933

reason, most of our analyses in Section 4 involve only relaxed haloes.

2.1.3 Shape parameter

Determining the shape of a three-dimensional distribution of particles is a non-trivial task (e.g. Jing & Suto 2002). Following Allgood et al. (2006), we determine the shapes of our haloes starting from the inertia tensor. As a first step, we compute the halo's 3×3 inertia tensor using all the particles within the virial radius. Next, we diagonalize the inertia tensor and rotate the particle distribution according to the eigenvectors. In this new frame (in which the moment of inertia tensor is diagonal) the ratios a_3/a_1 and a_3/a_2 (where $a_1 \geq a_2 \geq a_3$) are given by

$$\frac{a_3}{a_1} = \sqrt{\frac{\sum m_i z_i^2}{\sum m_i x_i^2}} \quad \frac{a_3}{a_2} = \sqrt{\frac{\sum m_i z_i^2}{\sum m_i y_i^2}}. \quad (4)$$

Next we again compute the inertia tensor, but this time only using the particles inside the ellipsoid defined by a_1 , a_2 and a_3 . When deforming the ellipsoidal volume of the halo, we keep the longest axis (a_1) equal to the original radius of the spherical volume (R_{vir}). We iterate this procedure until we converge to a stable set of axial ratios.

Since DM haloes tend to be more prolate on average (e.g. Kasun & Evrard 2005; Bett et al. 2007), a useful parameter that describes the shape of the halo is $q \equiv (a_2 + a_3)/2a_1$, with the limiting cases being a sphere ($q = 1$) and a needle ($q = 0$). A related parameter is the triaxiality $\tau \equiv (a_1^2 - a_2^2)/(a_1^2 - a_3^2)$, which can also be used to quantify the prolateness of haloes, but we find that it yields nearly identical results as the shape parameter, so we choose to use only q in Section 4.

Note that for haloes that are not spherical, the NFW fit to the density profile could be affected, potentially inducing an 'artificial' correlation between halo shape and the rms of the density profile fit. Nevertheless, although triaxial models are an improvement over spherical profile fits, Jing & Suto (2002) have shown that the differences between such fitted profiles are fairly small, and in any case, Macciò et al. (2007) have shown that there is no significant correlation between halo shape and the rms of the spherical profile fit.

2.1.4 Position angle Φ_L

We also consider the angle Φ_L between the major axis of the halo (a_1) and the angular momentum vector L . The angle Φ_L can be thought of as a proxy for the position of a possible stellar disc (and for the orientation of satellite galaxies) within the DM halo, assuming that the angular momentum of the disc and the one of the DM particles will be aligned (e.g. Sharma & Steinmetz 2005; Zentner et al. 2005b; Heller, Shlosman & Athanassoula 2007; Hahn, Teyssier & Carollo 2010). Nevertheless, it is difficult for simulations to reproduce the angular momentum distributions of observed disc galaxies (van den Bosch 2002; Sharma & Steinmetz 2005). More recently, Heller et al. (2007) claim to reproduce many of the properties of observed galactic discs, and the angular momentum of their simulated discs follows that of the DM and does not decrease too much while the discs form. Agertz et al. (2011) also can reproduce the angular momentum content of disc galaxies by assuming a lower star formation efficiency than other simulations and weaker feedback energy from supernovae. In contrast, Governato et al. (2010) argue that strong outflows from supernovae are

necessary to reproduce the angular momentum distribution of dwarf galaxies, and help to yield shallow DM profiles.

We note that the shape parameter q and Φ_L are not estimated from all of the particles in a halo, but rather from the halo's best-fitting ellipsoid. Nevertheless, for the vast majority of the haloes, these parameters are computed using >80 per cent of the particles, so we expect that this does not affect or bias our results significantly. In addition, observational estimates of the position angles of galaxies also have large uncertainties (e.g. Barnes & Sellwood 2003). In any case, although for a spherical halo Φ_L will be poorly defined, this does not result in systematic uncertainties in our analysis or in artificial correlations between Φ_L and halo shape.

2.1.5 Environment parameter Δ_8

Finally, in some of our analyses we also consider the overdensity (Δ_8), computed within a sphere of radius $R = 8 h^{-1}$ Mpc centred on the halo centre. Such an overdensity within a fixed aperture allows us to quantify the large-scale environment of haloes. Although most haloes are smaller than $2 \text{ Mpc } h^{-1}$ in size (e.g. Navarro et al. 1997), it is useful to focus on larger scales, in order to more clearly encompass non-linear scales within the aperture (e.g. Abbas & Sheth 2007).

We use the $8 \text{ Mpc } h^{-1}$ overdensity, Δ_8 , in our analyses in Section 4. In Appendix A, we show similar results for smaller scale overdensities, Δ_2 and Δ_4 .

2.2 Relaxed and unrelaxed haloes

Our halo finder, and halo finders in general, does not distinguish between relaxed and unrelaxed haloes (although there are recent exceptions, using halo phase-space densities: e.g. Zemp et al. 2009; Behroozi et al., in preparation). There are many reasons why we might want to remove unrelaxed haloes. First and foremost, unrelaxed haloes often have poorly defined centres, which makes the determination of a radial density profile, and hence of the concentration parameter, an ill-defined problem. Moreover, unrelaxed haloes often have shapes that are not adequately described by an ellipsoid, making our shape parameters ill defined as well.

Following Macciò et al. (2007, 2008), we decide to use a combination of two different parameters ρ_{rms} and x_{off} to determine the dynamical status of a given DM halo. The first quantity ρ_{rms} is defined as the rms of the NFW fit to the density profile (performed to compute c_{vir}). While it is true that ρ_{rms} is typically high for unrelaxed haloes, haloes with relatively few particles also have a high ρ_{rms} (due to Poisson noise) even when they are relaxed; furthermore, since the spherical averaging used to compute the density profiles has a smoothing effect, not all unrelaxed haloes have a high ρ_{rms} .

In order to circumvent these problems, we combine the value of ρ_{rms} with the x_{off} parameter, defined as the distance between the most bound particle (used as the centre for the density profile) and the centre of mass of the halo, in units of the virial radius. This offset is a measure for the extent to which the halo is relaxed: relaxed haloes in equilibrium will have a smooth, radially symmetric density distribution, and thus an offset that is virtually equal to zero. Unrelaxed haloes, such as those that have only recently experienced a major merger, are likely to reveal a strongly asymmetric mass distribution, and thus a relatively large x_{off} .

Although some unrelaxed haloes may have a small x_{off} , the advantage of this parameter over, for example, the actual virial ratio, $2T/V$, as a function of radius, is that the former is trivial to evaluate. In addition, the virial ratio can be quite noisy (e.g. Macciò, Murante & Bonometto 2003), while the substructure fraction,

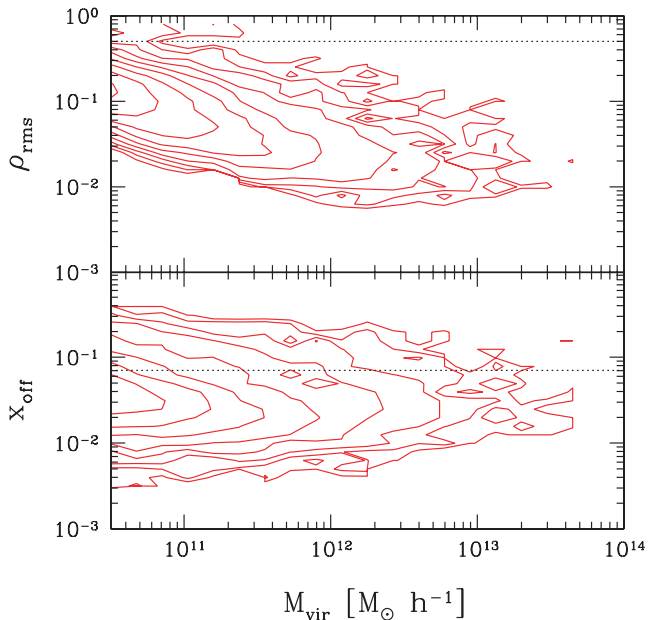


Figure 1. Correlation between the relaxedness parameters (x_{off} and ρ_{rms} and the virial mass M_{vir}). The two dotted lines represent the limit values used to define ‘good’ haloes.

another proxy for virialization, depends significantly on the simulation’s resolution (Macciò et al. 2007).

Fig. 1 shows the correlation between x_{off} (lower panel), ρ_{rms} (upper panel) and the virial mass of the halo for the W5-90.2 box. As already noted in Macciò et al. (2007), there is a correlation between the M_{vir} (or N_{vir}) and ρ_{rms} , which simply reflects the fact that for low halo masses the density profiles are more noisy. In most of our analyses in Section 4, we use a DM particle threshold as well as a mass criterion in order to minimize any effect of this correlation on our results. Finally, x_{off} and M_{vir} are extremely weakly correlated, if at all, and we find no correlation between x_{off} and ρ_{rms} .

Following Macciò et al. (2007), we split our halo sample into unrelaxed and relaxed haloes. The latter are defined as the haloes with $\rho_{\text{rms}} < 0.5$ and $x_{\text{off}} < 0.07$. About 70 per cent of the haloes in our sample qualify as relaxed haloes. In what follows, we will present results for two different samples of haloes: ‘all’, which includes all haloes with $N_{\text{vir}} > 1000$, and ‘good’ or ‘relaxed’, which is the corresponding subsample of relaxed haloes.

3 PRINCIPAL COMPONENT ANALYSIS

3.1 Introduction to PCA

PCA is a type of multivariate (multidimensional) statistics. The goal of such an analysis is to find a suitable representation of multivariate data, where representation here means that we somehow transform the data so that their essential structure is made more visible or accessible. PCA belongs to the area of unsupervised learning, because the representation must be learned from the data themselves without any external input from a supervising ‘teacher’.

PCA is a powerful technique that has been used in classification and dimensional reduction of large data sets, in order to decrease the complexity of the analysis. Each member of a data set (a DM halo, in our case) is defined by an ‘information vector’, an n -tuple of numbers (the halo’s relevant properties, in our case). PCA consists of de-correlating sets of vectors by performing rotations in

the n -dimensional parameter space. The final result is a diagonal covariance matrix, with the eigenvalues representing the amount of information (in the sense of variance) stored in each eigenvector, which is called a principal component (PC).

The PCs can be considered a set of basis vectors that optimally filter the information hidden in the data. If most of the variance in the original data can be accounted for by just a few of the components, we will have found a simpler description of the original data set. By convention, the first PC corresponds to the largest eigenvalue, the second PC corresponds to the second largest eigenvalue, and so on. A consequence of this definition is that the first PC is a minimum distance fit to a line in the space of the original variables. The second PC is then a minimum distance fit to a line in the plane perpendicular to the first PC, and so on. Therefore, the properties that dominate the first PC are the properties that contain the largest amount of information and to some extent determine the values of the other properties. Typically, the eigenvalue for a PC must be ≥ 1 to be significant; less significant PCs tend to be dominated by noise (e.g. Connolly & Szalay 1999). In any case, the amount of variance that a PC accounts for is more meaningful than the precise value of its eigenvalue. If the data have high signal-to-noise ratio, and if the data are not very incomplete, then the amount of variance of the first few PCs can be associated with ‘information’ contained in the data.

In astronomy, PCA applications in studies of multivariate distributions have been discussed in detail (Efstathiou & Fall 1984; Murtagh & Heck 1987). PCA methods have been used to study stellar, galaxy and quasar spectra (e.g. Connolly et al. 1995; Bailer-Jones, Irwin & von Hippel 1998; Madgwick et al. 2003; Yip et al. 2004; Ferreras et al. 2006; Rogers et al. 2007; Chen et al. 2009; Boroson & Lauer 2010; McGurk, Kimball & Ivezić 2010), galaxy properties (e.g. Conselice 2006; Scarlata et al. 2007) and stellar populations and the Fundamental Plane (e.g. Faber 1973; Eisenstein et al. 2003; Woo, Courteau & Dekel 2008). In a recent study, Bonoli & Pen (2009) analysed the stochasticity of DM haloes, which quantifies the scatter in the correlation between halo and galaxy density fields.

PCA methods have not yet been exploited to analyse the properties of DM haloes. Considering the increasing size and resolution of numerical simulations, PCA is a useful tool to determine the relative importance of halo properties, and to determine how fundamental the masses of haloes are. The primary goal of this paper is to use PCAs to analyse a variety of halo properties, including those introduced in Section 2.

3.2 Applying PCAs to DM haloes

Before we perform our PCAs on our simulation data, we subtract the mean from each halo property and normalize by the standard deviation. It follows that the distributions of the properties may be important: for example, choices about whether to take the logarithm of a property or whether to include objects whose properties are outliers in the distributions may be important. It is particularly useful to take the logarithm of halo mass, for example, because its range spans many orders of magnitude, and the spin parameter, because this makes its distribution more Gaussian; for consistency, however, we take the logarithms of all of the halo properties in our PCA. We also choose to exclude the extreme outliers ($> 5\sigma$) in the distributions of properties. In addition, some properties are so strongly correlated that one property almost completely determines another (such as the virial mass, radius and velocity, or spin and angular momentum, or shape and triaxiality), and in such a case we

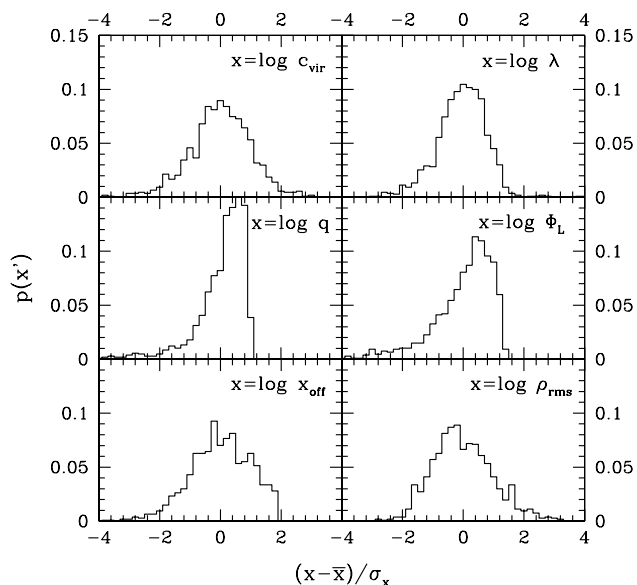


Figure 2. Distributions of halo properties at fixed mass [$\log(M/h^{-1} M_{\odot}) \sim 12$], for haloes with more than 1000 particles in the high-resolution 90 Mpc box. For a given property x , we define the mean-subtracted and standard deviation normalized quantity as $x' \equiv (x - \bar{x})/\sigma_x$.

Table 2. Median, mean and standard deviation of halo properties at fixed mass in the 90 Mpc box, whose distributions are shown in Fig. 2.

Property	Median	Mean	Std dev
$\log c_{\text{vir}}$	1.043	1.039	0.107
$\log \lambda$	-1.498	-1.513	0.252
$\log q$	-0.101	-0.126	0.094
$\log \Phi_L$	0.207	0.147	0.249
$\log x_{\text{off}}$	-1.614	-1.622	0.248
$\log \rho_{\text{rms}}$	-1.439	-1.413	0.259

include only what we believe is the more important property. By reducing the number of parameters, this tends to clarify the results of the PCA. We have performed numerous tests and have verified that these choices do not significantly affect our results.

We perform PCAs of each of the simulations described in Section 2. We analyse all of the haloes in each box, as well as the ‘relaxed’ haloes, with low values of ρ_{rms} (the rms of the NFW fit to the density profile) and x_{off} (the offset between the most bound particle and the centre of mass).

The probability distributions of the halo properties are shown in Fig. 2 and described in Table 2, for haloes in the 90 Mpc box, at fixed halo mass (see Section 4.2). By construction, PCA only finds independent normal processes, and works most effectively when the data set is jointly normally distributed. Our halo parameters are chosen to be approximately independent, and as can be seen in the figure, their distributions are approximately Gaussian. The shape parameter q and the position angle Φ_L have slightly non-Gaussian distributions, with a tail of low values, but this is not sufficient to significantly affect the PCA results.

Lastly, we show the log halo mass distributions in Fig. 3, for some of the analyses in Section 4. (In particular, these are the distributions for the PCAs whose results are shown in Tables 3, 5, 7 and 8.) We will perform PCAs with halo mass thresholds (Section 4.1),

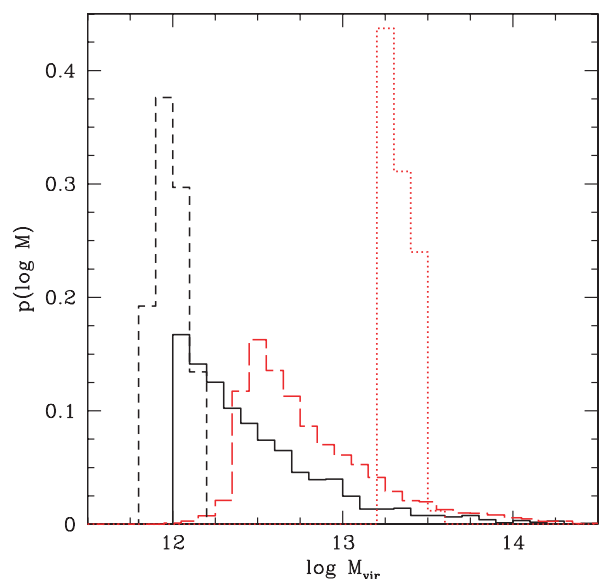


Figure 3. Log halo mass distributions for $\log(M/h^{-1} M_{\odot}) \geq 12$, 90 Mpc box (solid black histogram); $\log(M/h^{-1} M_{\odot}) \geq 12$, 180 Mpc box (long-dashed red histogram); relaxed haloes with $11.85 \leq \log(M/h^{-1} M_{\odot}) < 12.15$, 90 Mpc box (short-dashed black histogram); relaxed haloes with $13.2 \leq \log(M/h^{-1} M_{\odot}) < 13.5$, 180 Mpc box (dotted red histogram).

followed by PCAs at fixed halo mass (Section 4.2), and the mass distributions of some of these halo catalogues are shown in the figure. When a halo mass threshold is used, the mass distributions are highly non-Gaussian; none the less, the effect of this on the PCA results is minimal, as we discuss in the next section.

4 RESULTS

For the PCAs, we focus on eight halo properties: mass, concentration, spin, shape, Φ_L , x_{off} , ρ_{rms} and overdensity. First, in Section 4.1, we perform PCA of the haloes in our simulations using a halo mass threshold. Then in Section 4.2, we perform PCA at a fixed halo mass. Lastly, we briefly discuss the lack of a correlation between Φ_L and other halo parameters in Section 4.3.

4.1 PCAs with a halo mass threshold

The results of our PCAs of all of the haloes in the 90 and 180 Mpc boxes with $\log(M/h^{-1} M_{\odot}) \geq 12$ are shown in Tables 3 and 5; the results for the subset of relaxed haloes are shown in Tables 4

Table 3. All haloes with $\log(M/h^{-1} M_{\odot}) \geq 12$, 90 Mpc box, high resolution. $N = 2183$. Numbers in bold indicate dominant halo properties in a given principal component.

Property	PC1	PC2	PC3	PC4
$\log M_{\text{vir}}$	0.103	- 0.761	0.148	0.500
$\log c_{\text{vir}}$	- 0.555	0.156	-0.015	-0.134
$\log \lambda$	0.369	0.009	0.062	-0.333
$\log q$	0.094	0.579	0.143	0.759
$\log \Phi_L$	0.001	-0.059	- 0.971	0.181
$\log x_{\text{off}}$	0.547	-0.018	-0.012	-0.041
$\log \rho_{\text{rms}}$	0.488	0.240	-0.107	-0.105
Eigenvalues	2.56	1.18	1.01	0.90
Per cent of variance	36.5	16.9	14.4	12.8

Table 4. ‘Good’ (relaxed) haloes with $\log(M/h^{-1} M_{\odot}) \geq 12$, 90 Mpc box, high resolution. $N = 1701$.

Property	PC1	PC2	PC3	PC4
$\log M_{\text{vir}}$	0.163	− 0.762	0.064	0.166
$\log c_{\text{vir}}$	− 0.604	0.231	0.006	−0.114
$\log \lambda$	0.392	0.102	0.128	−0.221
$\log q$	0.085	0.307	0.304	0.887
$\log \Phi_L$	0.033	−0.025	− 0.928	0.300
$\log x_{\text{off}}$	0.543	−0.001	0.021	−0.047
$\log \rho_{\text{rms}}$	0.389	0.511	−0.162	−0.176
Eigenvalues	1.84	1.35	1.01	0.96
Per cent of variance	26.3	19.2	14.5	13.7

Table 5. All haloes with $\log(M/h^{-1} M_{\odot}) \geq 12$, 180 Mpc box. $N = 6391$.

Property	PC1	PC2	PC3	PC4
$\log M_{\text{vir}}$	0.480	0.467	0.025	−0.078
$\log c_{\text{vir}}$	0.221	− 0.639	−0.022	0.007
$\log \lambda$	−0.320	0.417	−0.033	0.046
$\log q$	−0.218	−0.023	0.263	− 0.936
$\log \Phi_L$	0.013	0.001	0.963	0.268
$\log x_{\text{off}}$	− 0.539	0.257	−0.040	0.167
$\log \rho_{\text{rms}}$	− 0.530	−0.358	−0.011	0.126
Eigenvalues	1.94	1.64	1.00	0.95
Per cent of variance	27.7	23.5	14.3	13.6

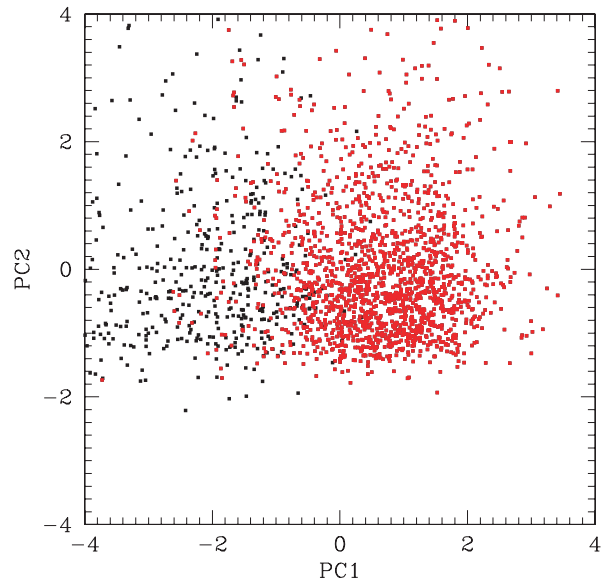
Table 6. ‘Good’ (relaxed) haloes with $\log(M/h^{-1} M_{\odot}) \geq 12$, 180 Mpc box. $N = 3757$.

Property	PC1	PC2	PC3	PC4
$\log M_{\text{vir}}$	− 0.618	−0.098	0.030	0.011
$\log c_{\text{vir}}$	0.443	− 0.469	−0.003	0.044
$\log \lambda$	−0.062	0.664	−0.044	−0.012
$\log q$	0.113	0.286	0.404	0.802
$\log \Phi_L$	0.010	−0.004	0.899	−0.435
$\log x_{\text{off}}$	0.236	0.494	−0.164	−0.406
$\log \rho_{\text{rms}}$	0.591	0.066	0.002	−0.008
Eigenvalues	2.06	1.25	1.01	0.98
Per cent of variance	29.4	17.8	14.4	14.0

and 6. Similar PCA results for smaller and larger simulation boxes are shown in Appendix A. Note that because of our resolution constraint, in the 180 Mpc box, most of the haloes with $\log(M/h^{-1} M_{\odot}) \geq 12$ are more massive than $10^{12.3} h^{-1} M_{\odot}$ (see Fig. 3).

We consider the following seven halo properties, of which we use the logarithms: the mass (M_{vir}), the concentration (c_{vir}), the spin (λ), the shape (q), the angle between the major axis and the total angular momentum (Φ_L), ρ_{rms} and x_{off} . We have performed tests of the uncertainties of the PCAs with bootstrap techniques and have found that our errors are extremely small; hence, we have omitted them from the tables, for clarity. In addition, our results are not significantly dependent on the selection criteria (i.e. the choice of the halo mass threshold, and for the PCAs later in this section, the particle limit and the x_{off} and ρ_{rms} criteria).

One striking result is the *lack* of a single or pair of dominant PCs. The first two PCs usually account for only about 40–50 per cent of the variance. In contrast, in PCAs of galaxy spectra, the first one or two PCs usually dominate (e.g. Connolly et al. 1995; Madgwick

**Figure 4.** First two PCs plotted against each other, for all haloes with $\log(M/h^{-1} M_{\odot}) \geq 12$ in the 90 Mpc box. We are using the results from the first two columns of Table 3: for example, for a given halo, $\text{PC1} = -0.103 (\log M_{\text{vir}})' + 0.555 (\log c_{\text{vir}})' - 0.369 (\log \lambda)' - 0.094 (\log q)' - 0.001 (\log \Phi_L)' - 0.547 (\log x_{\text{off}})' - 0.488 (\log \rho_{\text{rms}})'$, where the halo parameters are mean-subtracted and normalized by the standard deviation. The red points are the subset of ‘good’ (i.e. relaxed) haloes.

et al. 2003). Recall that, as stated in Section 3.1, the later PCs (with lower eigenvalues) tend to be dominated by noise.

For DM haloes, no parameter or combination of parameters strongly determines the other parameters. (This is partly by construction, because for closely related parameters, like mass and radius or shape and triaxiality, we have included only one of the parameters.) In other words, although halo properties are correlated with each other, we find no clearly fundamental correlation (nor a Fundamental ‘Plane’).

This is also evident in Fig. 4, in which the first two PCs of all haloes in the 90 Mpc box (Table 3) are shown. Some structure is apparent in the figure, with relaxed and unrelaxed haloes exhibiting different distributions. The figure illustrates that the relaxed haloes have larger values of PC1 (which is anticorrelated with the relaxedness parameters), and the fact that most haloes have $\text{PC2} > -2$ is simply due to the halo mass threshold.

The PCA results of the different simulations are similar, but not equivalent, most likely because of the different ranges of halo masses. In any case, in the majority of our simulation boxes, combinations of M_{vir} , c_{vir} , x_{off} and ρ_{rms} dominate the first two PCs. Halo mass and concentration, on the one hand, and halo relaxedness, on the other, can explain at least half of the variance. Note, however, that halo mass is never significant on the first PC. x_{off} and ρ_{rms} indicate the relaxedness of haloes, and for unrelaxed haloes, the concentration is not well defined. Therefore, *we argue that the relaxedness of haloes is at least as important as their mass and concentration* for determining halo properties. This is the primary result of our paper.

When halo parameters appear together on a PC, this can indicate a correlation between those parameters. For example, when a PC is dominated by mass and concentration, they usually have opposite signs, consistent with the fact that these halo properties are anticorrelated (e.g. Bullock et al. 2001). Correlated halo properties do not *necessarily* dominate the same PCs, however, and properties

dominating a PC are not necessarily correlated, although this is often the case. If two correlated properties have substantial scatter between them, or if their correlation is due to more fundamental correlations with a third parameter, or if two properties have strongly correlated errors, the interpretation of PCA results may be more complicated. With our choice of halo properties and selection criteria, we have attempted to minimize these complications.

Finally, note the similarities and differences between the results for all haloes and for just the relaxed haloes. Most of the same halo parameters dominate the first few PCs. In addition, the PCA results are mostly similar in the different resolution simulation boxes, indicating that our results are not very sensitive to resolution, though the significance of the spin parameter varies slightly. Although we find that the PCA results presented here are statistically stable, in some cases the first PC of one simulation box resembles the second PC of another, and vice versa, suggesting that it is the combination of the first two PCs that is important. In any case, throughout our analyses, halo mass and concentration, on the one hand, and relaxedness parameters (quantified by x_{off} and ρ_{rms}), on the other, continue to dominate the first two PCs.

We also note that the halo mass distributions are non-Gaussian when a mass threshold is used (Fig. 3), and this could potentially affect the PCA results. We have tested this effect by performing PCAs with subsamples that have Gaussian mass distributions, and by PCAs with the whole samples but with the mass distributions rescaled (forcing them to be Gaussian, with the same rank order), and we find that the results are virtually the same, differing from Tables 3–6 by <1 per cent.

None the less, some halo parameters are accurately determined only for relaxed haloes (Macciò et al. 2007), such that the very unrelaxed haloes introduce noise in the other parameters and could relegate them to the later PCs, or could introduce artificial correlations with the relaxedness parameters. For these reasons, for the rest of this paper, we focus on PCA results based on the subset of relaxed haloes.

4.2 PCAs at fixed halo mass

As shown in Fig. 2, at fixed halo mass, the distributions of the logarithm of most halo properties are approximately Gaussian [i.e. $p(\log \lambda|M)$, $p(\log c|M)$, etc., can be fitted fairly well by Gaussian distributions]. It therefore makes sense to perform PCAs of these variates at fixed mass. Constraints from these PCAs could be useful for halo model analyses (e.g. of galaxy clustering, weak lensing, galaxy groups and clusters), and perhaps for semi-analytic models as well, because they reduce the complexity of such analyses to a Gaussian mixture model. Gaussian mixture models have a well-developed statistical framework and have been applied to a variety of astrophysical studies (e.g. Kelly & McKay 2004; Hao et al. 2009; Skibba & Sheth 2009).

We first perform PCAs similar to those in Section 4.1, but now at fixed halo mass, using the other six halo parameters. The results of these PCAs are shown in Table 7 for haloes with $11.85 \leq \log(M/h^{-1} M_{\odot}) < 12.15$ in the 90 Mpc box, and in Table 8 for haloes with $13.2 \leq \log(M/h^{-1} M_{\odot}) < 13.5$ in the 180 Mpc box. Only haloes with more than 1000 particles are included.

The results are similar to those of the previous section, with concentration and x_{off} dominating the first PC, although the spin parameter is now important on this PC as well. In addition, Φ_L and the shape parameter dominate the second PC. Halo model analyses often involve constructing mock galaxy catalogues, and the models are often based on the clustering and abundance of haloes as a

Table 7. ‘Good’ (relaxed) haloes with $\log(M/h^{-1} M_{\odot}) \sim 12$ and more than 1000 particles, 90 Mpc box, high resolution. $N = 1056$.

Property	PC1	PC2	PC3	PC4
$\log c_{\text{vir}}$	– 0.601	–0.076	0.044	–0.011
$\log \lambda$	0.441	–0.089	0.544	–0.048
$\log q$	0.167	– 0.591	– 0.427	– 0.654
$\log \Phi_L$	–0.023	0.772	–0.113	– 0.605
$\log x_{\text{off}}$	0.552	0.088	0.158	–0.071
$\log \rho_{\text{rms}}$	0.334	0.183	– 0.694	0.446
Eigenvalues	1.79	1.06	0.97	0.93
Per cent of variance	29.8	17.6	16.2	15.4

Table 8. ‘Good’ (relaxed) haloes with $\log(M/h^{-1} M_{\odot}) \sim 13.35$ and more than 1000 particles, 180 Mpc box. $N = 416$.

Property	PC1	PC2	PC3	PC4
$\log c_{\text{vir}}$	– 0.643	0.107	0.015	0.000
$\log \lambda$	0.496	0.141	0.339	0.120
$\log q$	0.181	– 0.567	–0.389	– 0.642
$\log \Phi_L$	0.052	– 0.588	0.715	0.058
$\log x_{\text{off}}$	0.520	0.058	–0.353	0.344
$\log \rho_{\text{rms}}$	–0.186	– 0.546	–0.313	0.672
Eigenvalues	1.55	1.13	0.97	0.93
Per cent of variance	25.8	18.8	16.2	15.6

Table 9. ‘Good’ (relaxed) haloes with $\log(M/h^{-1} M_{\odot}) \sim 13.3$ and more than 1000 particles, 180 Mpc box. $N = 383$.

Property	PC1	PC2	PC3	PC4
$\log \Delta_8$	–0.252	–0.103	0.761	–0.036
$\log c_{\text{vir}}$	0.630	0.059	0.247	0.049
$\log \lambda$	– 0.499	–0.054	0.123	0.226
$\log q$	–0.126	– 0.585	0.115	– 0.667
$\log \Phi_L$	–0.007	– 0.595	– 0.495	0.159
$\log x_{\text{off}}$	0.509	0.205	–0.149	0.187
$\log \rho_{\text{rms}}$	0.128	– 0.494	0.256	0.664
Eigenvalues	1.51	1.10	1.02	0.97
Per cent of variance	21.6	15.7	14.6	13.8

function of their masses; at fixed mass, the models use halo concentrations and density profiles. Perhaps the models could also use halo spin distributions, to more accurately describe galaxy velocities, as well as x_{off} distributions, to describe central galaxies offset from the halo centres.

Next, we perform PCAs with an additional parameter, the overdensity Δ_8 within a sphere of radius 8 Mpc, which is an indicator of a halo’s local environment (since the largest haloes have virial radii of ~ 2 Mpc/h). At a fixed halo mass, there may be a range of environments: some haloes may be embedded within large structures or filaments or surrounded by infalling haloes, while some may be relatively isolated (e.g. Colberg et al. 1999). Using a fixed halo mass also ensures that the haloes are of similar size and are much smaller than the density aperture. We have restricted this analysis to haloes with more than 1000 particles and with $13.15 < \log(M/h^{-1} M_{\odot}) < 13.45$ in the 180 Mpc box and $13.45 < \log(M/h^{-1} M_{\odot}) < 13.75$ in one of the 300 Mpc boxes. It is useful to analyse the haloes in a large simulation box, in which cosmic variance is relatively unimportant. The results of these PCAs are shown in Tables 9 and 10.

Table 10. ‘Good’ (relaxed) haloes with $\log(M/h^{-1} M_{\odot}) \sim 13.6$ and more than 1000 particles, 300 Mpc box. $N = 944$.

Property	PC1	PC2	PC3	PC4
$\log \Delta_8$	-0.138	-0.164	0.819	0.418
$\log c_{\text{vir}}$	0.629	0.051	0.112	-0.061
$\log \lambda$	- 0.521	-0.050	0.095	0.138
$\log q$	-0.050	0.603	0.249	-0.149
$\log \Phi_L$	0.093	0.405	-0.347	0.840
$\log x_{\text{off}}$	- 0.544	0.256	-0.206	-0.183
$\log \rho_{\text{rms}}$	0.084	0.613	0.288	-0.203
Eigenvalues	1.50	1.11	1.02	0.96
Per cent of variance	21.5	15.9	14.6	13.8

We find that the local overdensity is usually significant on the third PC, where it is sometimes paired with the angle Φ_L . This is consistent with Altay, Colberg & Croft (2006), who found that haloes are aligned with their large-scale structures (if they are located in filaments, for example), while halo shapes are not dependent on membership of a filament. It is also consistent with Wang et al. (2011), who showed that a halo’s spin is correlated with the local overdensity and the strength of the tidal field. In any case, in all of our simulations, the density is rarely significant on the first two PCs. Therefore, we argue that, at fixed mass, a halo’s environment is relatively unimportant, because it does not significantly determine the halo’s properties. At fixed mass, a halo’s angular momentum and degree of relaxedness are much more important.

4.3 Correlation with Φ_L

As previously stated, the angle Φ_L can be thought of as a proxy for the position of a possible stellar disc within the DM halo, assuming an alignment between the angular momentum of the DM particle and the stellar component (e.g. Hahn et al. 2010). Fig. 5 shows the

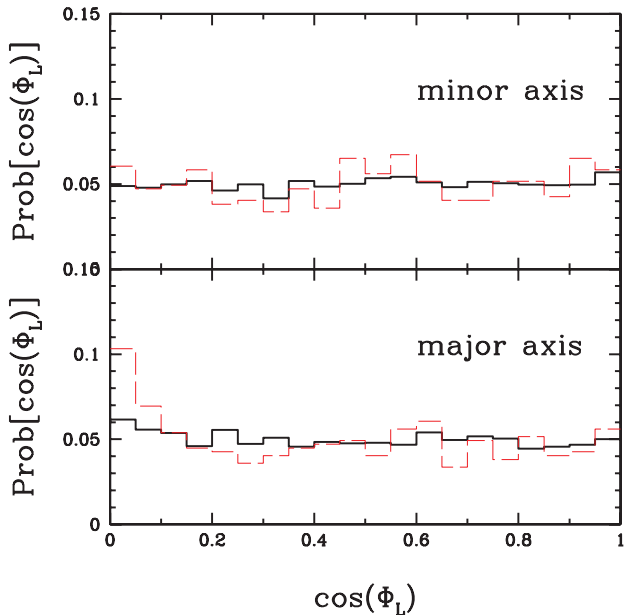


Figure 5. Probability distribution for the angle Φ_L between the angular momentum vector L and the minor (upper panel) and major (lower panel) axis of the DM distribution. The solid black line is for relaxed haloes in the 90 Mpc box, while the red line is for the spin haloes ($\lambda > 0.07$, ≈ 12 per cent of the total).

probability distribution of the cosine of Φ_L in the high-resolution 90 Mpc box. In the upper panel, the angle Φ_L is defined as the angle between the minor axis and the angular momentum L , while in the lower panel the major axis is used. For relaxed haloes the probability distribution is remarkably flat in both cases, and there is no sign of a possible alignment between either the major or minor axis and L . If we restrict our analysis to high spin haloes ($\lambda > 0.07$, red dashed line, ≈ 12 per cent of the total), a weak maximum at $\cos(\Phi_L) = 0$ arises for the angle with the major axis, while there is still no evidence for a correlation between the orientation of the minor axis and the halo angular momentum.

This result is consistent with Zentner et al. (2005b) and Kuhlen, Diemand & Madau (2007), who showed that subhaloes are distributed anisotropically and are preferentially aligned with the major axis of the triaxial halo. Studies of satellite galaxies in groups have reached similar conclusions (Yang et al. 2006; Faltenbacher et al. 2007; Kang et al. 2007; Wang et al. 2008). None the less, other studies have detected significant correlations between the spin axis of haloes and their *minor* axis, stronger than correlations with the major axis (Bailin & Steinmetz 2005; Bett et al. 2007). It is possible that the disagreement between our result (Fig. 5) and these studies somehow owes to the differences between SO and friends-of-friends (FOF) halo-finding algorithms. As noted by a number of authors, there are different disadvantages to these two types of halo finders (e.g. Bett et al. 2007; Tinker et al. 2008b; Knebe et al. 2011). In particular, SO halo finders tend to impose a more spherical geometry on the resulting systems. On the other hand, FOF halo finders often identify objects linked with neighbouring objects via tenuous bridges of particles, resulting in anomalously large velocity dispersions, masses and spin parameters of ≈ 6 per cent of objects (Bett et al. 2007). In addition, FOF halo finders define a halo centre as the centre of mass, while most other algorithms (including ours) define the centre as the centre of the density profile (i.e. the most bound particle; Knebe et al. 2011). These differences with FOF halo finders may explain the disagreement about correlations between the spin axis and major or minor axes of haloes.

Returning to our PCA results, in the different simulation boxes, rarely does a single halo parameter dominate a particular PC. None the less, Φ_L usually dominates the third or fourth PC, although sometimes the shape parameter is also significant on the same PC. In other words, one might interpret this as evidence that Φ_L is weakly correlated with spin and/or shape, but not related to other halo properties. Φ_L , which is the angle between the major axis of the halo and its angular momentum vector, is expected to be important for the formation of disc galaxies, because the angular momenta of the stellar disc and the DM particles are generally, though not always, expected to be aligned (van den Bosch et al. 2002; Kazantzidis et al. 2004; Sharma & Steinmetz 2005; Zentner et al. 2005b; Bett et al. 2010; Hahn et al. 2010).

In contrast to this, we show in Fig. 6 that Φ_L is uncorrelated with the shape parameter q . This result is independent of the value of the angular momentum (quantified by the spin parameter) as shown in Fig. 7. Therefore, as a consequence of our results, we argue that it is not possible to reliably determine the angle of the angular momentum vector from a halo’s shape. This has consequences for the formation of disc galaxies, which we discuss in the final section of the paper.

5 DEPENDENCE ON HALO RELAXEDNESS

From the results of the PCAs in Sections 4.1 and 4.2, we argued that halo ‘relaxedness’ is as important as halo mass in determining halo

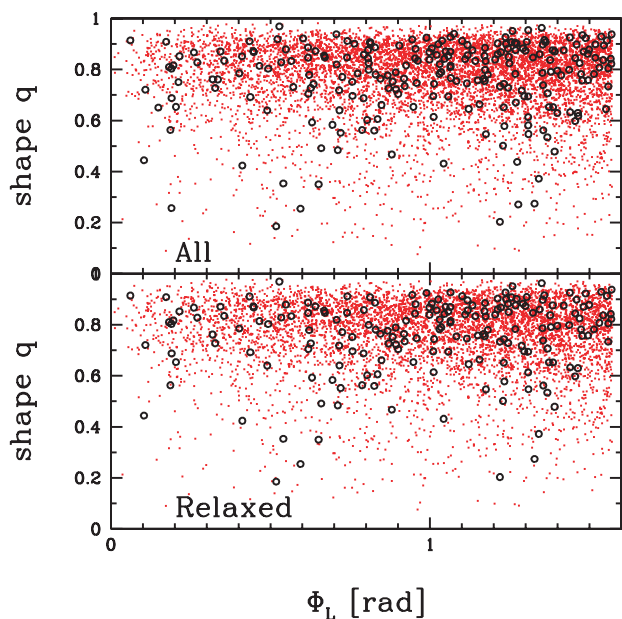


Figure 6. Relation between the shape (q) and the angle Φ_L . Black circle and red dots represent haloes from the W5-30.1 and the W5-90.2 simulations, respectively. The two panels show results for all (upper) and relaxed haloes. In both cases there is no correlation between the two properties.

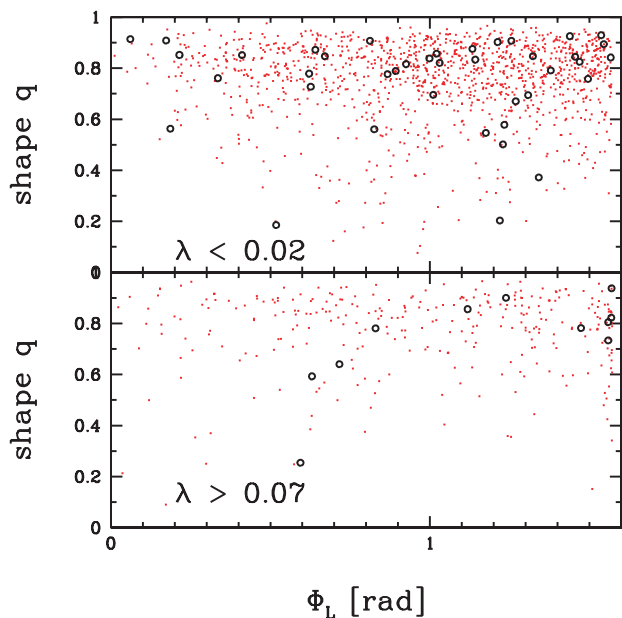


Figure 7. Same as Fig. 6 but of haloes with different spins. The upper panel shows results for relaxed haloes with little angular momentum (defined as haloes with $\lambda < 0.02$), while the lower one is for relaxed haloes with a lot of angular momentum ($\lambda > 0.07$). Again, there is no correlation between spin (λ) and the angle Φ_L .

properties, because mass and relaxedness account for the majority of the variance in the PCAs. This conclusion about the importance of halo relaxedness is consistent with other studies, such as Shaw et al. (2006), who find that the dynamical state of haloes, quantified by the virial ratio, is correlated with halo parameters: mass, concentration, spin as well as the substructure fraction. In addition, Neto et al. (2007) define an ‘equilibrium state’ in terms of the substructure fraction, centre-of-mass displacement (like our x_{off}) and virial ratio,

and find that a halo’s equilibrium state is correlated with its mass, concentration and spin.

In order to investigate this issue further, in this section we quantify and compare correlations between halo properties and mass, and between them and the relaxedness parameters.

In particular, we first perform Spearman rank tests¹ on some correlations between these parameters. We begin with correlations with halo relaxedness parameters at fixed mass [$11.85 \leq \log(M/h^{-1} M_{\odot}) < 12.15$], using haloes with more than 1000 particles in the 90 Mpc box. This is one of the catalogues used in Section 4.2 (see PCA result in Table 7). The correlations between concentration and x_{off} , and between spin and x_{off} are statistically significant, with Spearman ranks $r_s = -0.65$ and 0.42 , respectively, and remain significant if only ‘good’ (i.e. relaxed) haloes are selected. The correlation with shape is weaker, while that with Φ_L is not significant at all. Of the correlations with ρ_{rms} , only that with concentration is significant, with $r_s = -0.45$. Therefore, the dependence on relaxedness at fixed mass is predominantly due to x_{off} .

In order to perform a fair comparison with the halo mass dependence, we test the strength of correlations with mass at fixed relaxedness: we use the same simulation box and select haloes within a narrow range of relaxedness parameters ($x_{\text{off}} < 0.018$ and $\rho_{\text{rms}} < 0.14$), such that there is a similar number of haloes (~ 1300) as the previous test at fixed mass. Only the resulting mass–concentration correlation is still significant, with $r_s = -0.38$. These results imply that the expected weak correlations between mass and the spin parameter, and between mass and shape (e.g. Bett et al. 2007; Macciò et al. 2007), are no longer significant once one accounts for the dependence on halo relaxedness.

In Fig. 8, we show some of these correlations that we tested with the Spearman rank coefficients. In the left-hand panels of Fig. 8, we show the correlation between concentration and x_{off} at fixed mass, and compare it to the mass–concentration relation at fixed relaxedness parameters [similar to fig. 2 of Macciò et al. (2008)]. Halo mass and the centre-of-mass offset appear to be independently correlated with concentration (i.e. one is not due to the other). More massive haloes and haloes with larger offsets tend to be less concentrated.

We show the analogous plots with the spin parameter λ in the right-hand panels of Fig. 8. One can see no mass dependence at fixed relaxedness, although at fixed mass, haloes with larger offsets x_{off} tend to have larger spin. Although the spin parameter is not well defined for unrelaxed haloes, we see that the correlation occurs for haloes defined as relaxed [with $\log x_{\text{off}} < -1.15$, or the stricter criterion in Macciò et al. (2007), $\log x_{\text{off}} < -1.40$].

These tests show that halo concentration and spin are as tightly correlated with the relaxedness parameters (especially x_{off}) as with halo mass. This confirms our conclusion that the relaxed state of a halo is as important as its mass, in determining its other properties.

6 DISCUSSION AND CONCLUSIONS

We used numerical simulations of DM haloes to study a variety of halo properties: mass, concentration, spin, shape, the angle between the major axis and the angular momentum vector (Φ_L), the distance between the most bound particle and the centre of mass (x_{off}), the rms of the NFW fit to the density profile (ρ_{rms}), and the local

¹ The Spearman rank correlation coefficient may have a value between -1 and 1 . A positive (negative) value indicates an (anti)correlation, and a value of 0 indicates no correlation.

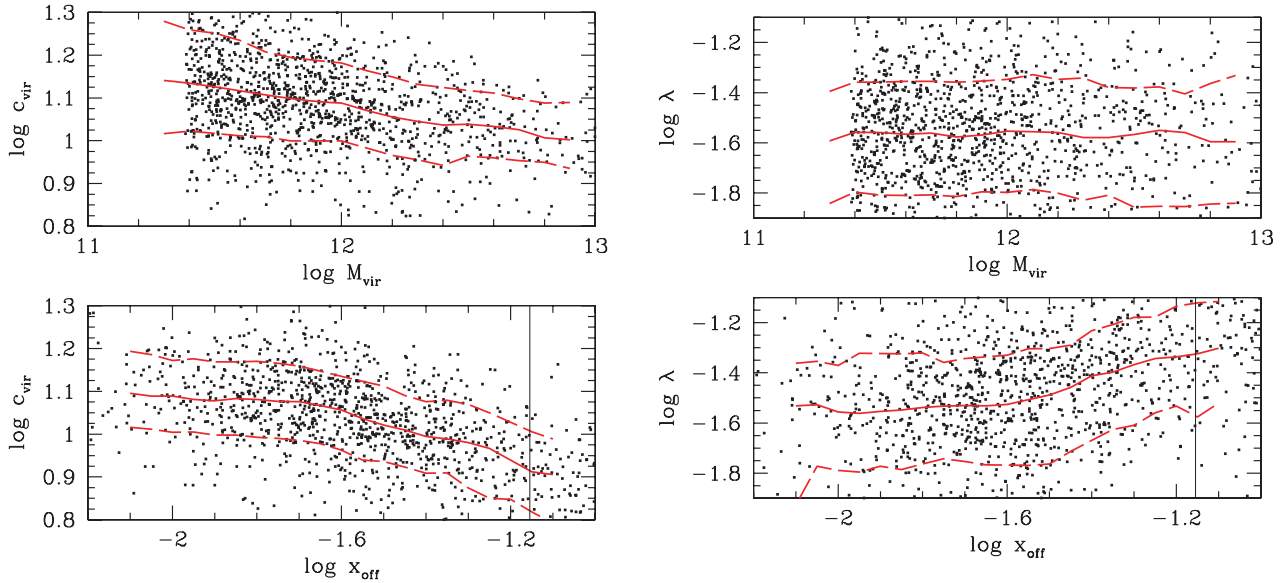


Figure 8. Left: M_{vir} versus c_{vir} for very relaxed haloes ($x_{\text{off}} < 0.018$, $\rho_{\text{rms}} < 0.14$) in the upper panel, and x_{off} versus c_{vir} at fixed mass ($11.85 \leq \log(M/h^{-1} M_{\odot}) < 12.15$, as in Section 4.2) in the lower panel, for haloes with more than 1000 particles in the 90 Mpc box. Right: same, but with the spin parameter, λ . Points indicate individual haloes, solid lines show the running medians, and dashed lines show the 16th and 84th percentiles. The vertical lines in the lower panels indicate our cut for relaxed haloes, at $x_{\text{off}} < 0.07$.

overdensity. We analysed all of these properties together, by employing PCA.

We summarize our main results as follows.

(i) There is no single dominant PC or pair of PCs. Unlike the spectra of galaxies, for DM haloes, there is no parameter or combination of parameters that strongly determines the other parameters.

(ii) We find that whether a halo is relaxed or not is at least as important as the halo’s mass and concentration. Even for relatively relaxed haloes, the degree of relaxedness, quantified by x_{off} and ρ_{rms} , is still as important as mass and concentration, and these four properties tend to dominate the first two PCs.

(iii) Φ_L is not significantly correlated with any other halo properties. It is therefore not possible to reliably estimate the direction of the angular momentum vector of a halo from its shape.

(iv) A halo’s ‘environment’ (quantified by the 8 Mpc overdensity) usually dominates the third PC, where it is sometimes combined with the angle Φ_L , but it is less significant than other halo properties on the first two PCs. Therefore, at fixed halo mass, the environment is relatively unimportant in determining the other halo properties.

Our results have many important implications, especially for halo occupation models and semi-analytic models. Most such models explicitly assume that halo mass (or the combination of mass and concentration) largely determines other important halo properties, and to some extent determines the properties of the galaxies hosted by a halo; however, our results show that this assumption is often false: halo mass is less dominant than one might expect. *The degree of relaxedness of a halo is at least as important as its mass.* Therefore, the relaxedness of a halo could affect the formation of galaxies hosted by it, independent of halo mass. In an unrelaxed halo, for example, galaxies may experience stronger tidal forces or their supply of gas to form stars could be disrupted. For semi-analytic models, one solution may be to quantify the relaxedness of the DM haloes and to use them in addition to the halo merger histories. For halo occupation models, it is possible that the unrelaxed-

ness of some haloes may affect or blur the relation between halo mass and central galaxy luminosity or stellar mass. In addition, it is possible that the halo mass function and halo mass–concentration relation assumed in the models could be different for relaxed and unrelaxed haloes. It is not yet clear how strongly halo relaxedness may affect results from halo occupation and semi-analytic models, and these issues deserve further study.

A related issue is that of halo ‘assembly bias’, in which the formation time of haloes (or other properties related to halo assembly) is correlated with the environment at fixed mass (e.g. Sheth & Tormen 2004; Croton et al. 2007; Wetzel et al. 2007; Fakhouri & Ma 2010; Faltenbacher & White 2010). The relaxedness of haloes, which we quantify with ρ_{rms} and x_{off} , is clearly related to their assembly: a halo that has recently formed or recently experienced a merger is more likely to be unrelaxed. Although our density parameter, Δ_8 , is subdominant in our PCAs, it is none the less possible that the relaxedness of haloes is related to their larger scale environment. Hence, the correlations with halo relaxedness may be a manifestation of halo assembly bias. None the less, assembly bias may have a relatively weak effect for halo models of galaxy clustering (e.g. Blanton & Berlind 2007; Tinker et al. 2008a; Skibba & Sheth 2009).

Our result for the angle Φ_L also has interesting implications. The fact that Φ_L is not significantly correlated with other halo properties may present a difficulty for some models. While the angular momentum of the stellar disc of a galaxy is thought to be aligned with the DM particles of the host halo (such that the pole of the disc is collinear with the angular momentum vector of the halo), *our result implies that it is actually quite difficult, if not impossible, to accurately determine the orientation of the disc simply by looking at the halo’s shape.* This is consistent with hydrodynamical simulations that include baryonic cooling, which result in discs with spin axes that are very poorly aligned with the halo (Bailin et al. 2005). It suggests that the direction of the angular momentum of the baryons is not well conserved throughout the disc formation process (van den Bosch et al. 2002; Yang et al. 2006). Heller et al. (2007) also argue that the shape and morphology of a

disc depend on the relative angle with the major axis of the halo, but this dependence too is not well conserved. Therefore, although the formation of disc galaxies and their host haloes are certainly related, after they have evolved, it is not clear that one could reliably determine to what extent the disc remains aligned with the halo.

Finally, at the time of publication, Jeesson-Daniel et al. (2011) posted a preprint of a related study, involving PCAs of various halo properties. Some of the results in Jeesson-Daniel et al. are consistent with ours, such as that halo mass, concentration, substructure and relaxedness are among a set of parameters that together account for a large fraction of the variance. None the less, their relaxedness parameter, which is similar to our x_{off} , accounts for less variance in their PCAs than in ours.

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APPENDIX A: SUPPLEMENTARY PCAS AND TESTS

In this appendix, we include additional PC analyses that can be compared to the results in Section 4.

First, we perform PCAs for halo mass threshold [$\log(M/h^{-1} M_{\odot}) > 12$] for the 40 and 300 Mpc simulation boxes, analogous to the results shown in Tables 4 and 6. The results of these PCAs are shown in Tables A1 and A2. The PCAs of the 40 and 90 Mpc boxes are similar, although for the 40 Mpc box, λ is now very significant on the second PC. The PCAs of the 180 and 300 Mpc boxes are similar as well, although for the 300 Mpc box, c_{vir} is more significant on the first PC than on the second one. In any case, even when one compares the PCAs of the 40 and 300 Mpc boxes, which have very different resolutions, the results are quite similar. We conclude that our results are not very sensitive to the resolution of the simulations.

Table A1. Relaxed haloes with $\log(M/h^{-1} M_{\odot}) \geq 12$, 40 Mpc box. $N = 1822$.

Property	PC1	PC2	PC3	PC4
$\log M_{\text{vir}}$	-0.618	0.078	-0.002	0.015
$\log c_{\text{vir}}$	0.408	0.499	-0.015	0.078
$\log \lambda$	-0.068	-0.645	-0.046	-0.070
$\log q$	0.093	-0.267	0.458	0.820
$\log \Phi_L$	0.002	-0.053	-0.881	0.467
$\log x_{\text{off}}$	0.292	-0.503	-0.106	-0.313
$\log \rho_{\text{rms}}$	0.595	-0.046	-0.014	-0.022
Eigenvalues	2.14	1.26	1.01	0.97
Per cent of variance	30.6	18.0	14.4	13.9

Table A2. Relaxed haloes with $\log(M/h^{-1} M_{\odot}) \geq 12$ and more than 1000 particles, 300 Mpc box. $N = 1905$.

Property	PC1	PC2	PC3	PC4
$\log M_{\text{vir}}$	-0.460	0.472	-0.124	-0.003
$\log c_{\text{vir}}$	0.609	0.185	0.045	-0.009
$\log \lambda$	-0.338	-0.465	0.116	0.044
$\log q$	0.075	-0.367	-0.360	-0.818
$\log \Phi_L$	0.112	-0.161	-0.862	0.463
$\log x_{\text{off}}$	-0.312	-0.483	0.136	0.240
$\log \rho_{\text{rms}}$	0.434	-0.364	0.279	0.239
Eigenvalues	1.62	1.33	0.99	0.96
Per cent of variance	23.2	19.0	14.2	13.7

Table A3. Relaxed haloes with $\log(M/h^{-1} M_{\odot}) \sim 13.3$ and more than 1000 particles, 180 Mpc box. $N = 383$.

Property	PC1	PC2	PC3	PC4
$\log \Delta_2$	0.282	0.406	0.559	0.031
$\log c_{\text{vir}}$	-0.621	0.221	0.150	-0.023
$\log \lambda$	0.511	0.035	0.181	-0.290
$\log q$	0.098	-0.542	0.201	0.755
$\log \Phi_L$	-0.023	-0.688	0.084	-0.580
$\log x_{\text{off}}$	0.505	0.061	-0.201	0.062
$\log \rho_{\text{rms}}$	-0.098	-0.118	0.738	-0.070
Eigenvalues	1.52	1.11	1.08	0.97
Per cent of variance	21.8	15.8	15.4	13.8

Table A4. ‘Good’ (relaxed) haloes with $\log(M/h^{-1} M_{\odot}) \sim 13.3$ and more than 1000 particles, 180 Mpc box. $N = 383$.

Property	PC1	PC2	PC3	PC4
$\log \Delta_4$	0.233	-0.181	0.751	0.049
$\log c_{\text{vir}}$	-0.626	0.024	0.303	-0.024
$\log \lambda$	0.502	-0.066	0.157	-0.271
$\log q$	0.124	-0.569	-0.110	0.772
$\log \Phi_L$	0.131	-0.572	-0.442	-0.470
$\log x_{\text{off}}$	0.521	0.200	-0.006	-0.037
$\log \rho_{\text{rms}}$	-0.123	-0.521	0.336	-0.325
Eigenvalues	1.51	1.10	1.03	0.97
Per cent of variance	21.6	15.7	14.7	13.8

Table A5. ‘Good’ (relaxed) haloes with $\log(M/h^{-1} M_{\odot}) \sim 13.3$ and more than 1000 particles, 180 Mpc box. $N = 383$. Like Table 9, but with 1σ error added to $\log \lambda$.

Property	PC1	PC2	PC3	PC4
$\log \Delta_8$	0.249	0.100	0.762	0.210
$\log c_{\text{vir}}$	-0.663	0.022	0.195	0.176
$\log \lambda$	0.282	0.367	-0.020	0.639
$\log q$	0.186	0.428	0.251	-0.713
$\log \Phi_L$	0.034	0.623	-0.504	0.034
$\log x_{\text{off}}$	0.602	-0.245	-0.084	0.043
$\log \rho_{\text{rms}}$	-0.143	0.472	0.240	0.073
Eigenvalues	1.38	1.12	1.02	1.00
Per cent of variance	19.7	16.0	14.5	14.3

Next, we perform more PCAs at fixed mass, using smaller scale overdensities Δ_2 and Δ_4 , in 2 and 4 Mpc apertures, analogous to the results in Table 9. The results of these PCAs are shown in Tables A3 and A4. When using Δ_4 , the first three PCs are essentially the same as when Δ_8 was used, and q still dominates the fourth PC. When Δ_2 is used, the second and third PCs are different, although they are dominated by the same parameters. In any case, regardless of the scale of the overdensity, we conclude that halo environment is less prominent in our PCAs, and therefore less important, than the other halo parameters.

Finally, while we have attempted to minimize the effects of correlated errors, we have not heretofore accounted for parameters with particularly large errors. One way to test the effects of such errors is to add error (with a Gaussian distribution) to a particular parameter. Such tests yield PCA results similar to those shown in Section 4, although we find that the first few PCs are usually more robust than the later PCs. An example is shown in Table A5, in which

we have significantly increased the error of the spin parameter, by 1σ . Comparing to Table 9, the numbers are slightly different, but the dominant parameters on each PC remain the same. The only exception is λ itself, which is no longer significant on PC1, but is instead slightly significant on PC2 and more significant on PC4. Similar tests with other parameters have yielded smaller changes than shown here: for example, adding error to the concentration does not affect any of the dominant parameters on the PCs, and

only slightly decreases the significance of c_{vir} on PC1. We have also tested the effects of asymmetric errors, but these significantly affect the PCA results only if they are larger and much more asymmetric than the expected errors in these simulations. We conclude that our PCA results are robust to the parameter errors.

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